

Giant oxygen isotope effect in high- T_c metal-oxide superconductors

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The theory of polaron superconductivity is shown to admit large values of the isotope effect, $\alpha > 0.5$. The oxygen isotope effect in $\text{YBa}_2\text{Cu}_3\text{O}_{7-x}$ can be anomalously large.

The first experiments¹ on the isotopic replacement of oxygen in the recently discovered high- T_c superconductors (YBCO and LSCO) showed that the value of the isotopic constant is relatively low:

$$\alpha = - \frac{d \ln T_c}{d \ln M} \lesssim 0.2. \quad (1)$$

This result was interpreted in several studies as evidence of the presence of soft modes,² anharmonicity, and Coulomb interaction or as an indication of the dominant role played by the nonphonon carrier-pairing mechanism (see, e.g., Ref. 3). Here M is the mass of the oxygen atom.

On the other hand, the energy-band calculations⁴ and analytic calculations,⁵ a

comparison of the tunnel characteristics with the neutron vibration spectra,⁶ the temperature dependence of thermal conductivity,⁷ and several other experiments⁸ have shown that the carriers of the high- T_c metal-oxide superconductors interact strongly with the clearly identifiable hard vibration modes of the light oxygen anions and that this interaction, in principle, should lead to a more pronounced isotope effect.

The low values of α in the experiments¹ on the replacement of O^{16} by O^{18} may stem from the fact that O^{16} does not leave the crystallographic positions in the copper planes and the primary candidate for a replacement is the more mobile oxygen in the copper chains and in the Ba neighborhood, whose isotopic state affects T_c only slightly. These conclusions have been confirmed by the recently obtained results⁹ on the synthesis of YBCO from the basic oxides which generally do not contain an O^{16} isotope:

$$\alpha \approx 2.5. \quad (2)$$

This result corresponds to a decrease in T_c from 92 K in O^{16} samples to 59 K and 77 K in samples which were synthesized identically from O^{18} and O^{17} isotopes, respectively. The giant isotope effect (2) is five times larger than the limiting value predicted by the BCS theory for a weak and strong electron-phonon coupling.

In this letter we offer an explanation of the anomalous isotope effect⁹ on the basis of a polaron model of a high-temperature superconductor,² according to which a sufficiently strong electron-phonon interaction ($\lambda \geq 1$) leads to the formation of small heavy bipolarons with an effective mass m^{**} which condense to a charged superfluid Bose liquid at the temperature¹⁰

$$T_c = f(p) / m^{**}, \quad (3)$$

where $f(p)$ is the concentration function of the carriers p , which does not depend on the mass of the isotope. At low concentrations ($\hbar = k_B = 1$) we have

$$f(p) \approx 3.3 (p/2)^{2/3}. \quad (4)$$

According to Aleksandrov and Kabanov,¹¹ in the simple case of an interaction with a dispersion-free vibration mode with a frequency ω we have

$$m / m^{**} = 2T_{ij} \Delta^{-1} M (1.1 + \Delta/\omega, 2g^2), \quad (5)$$

where $\Delta = (2g^2\omega - V_c)$ is the bipolaron binding energy which does not depend on the isotope mass, V_c is the Coulomb repulsion potential, g is the Fröhlich dimensionless coupling constant, T_{ij} is the integral which describes the hopping in the rigid lattice and which determines the band mass m , and $M(a, c, z)$ is a confluent hypergeometric function.

Assuming $\omega \sim M^{-1/2}$ and $g^2 \sim M^{1/2}$, we find from (3)–(5)

$$\alpha = g^2 F(\Delta/\omega, 2g^2), \quad (6)$$

where

$$F(x, y) = 1 + M^{-1} (1, 1 + x, y) (M(1, 2 + x, y) - (x/y) dM(1, 1 + x, y)/dx)$$

varies from $F(0, y) = 1$ to $F(\infty, y) = 2$.

As we can see from expression (6), the bipolaron superconductivity theory¹⁰ thus gives us the conceptual possibility of producing an isotope effect $\alpha > 0.5$, with values of $g^2 \gtrsim 1$, which account for the existence of a small-radius polaron.¹⁰ The value (2) corresponds to $g^2 = 1.25-2.5$, in agreement with the estimates² of the bipolaron effective mass obtained from T_c , the magnetic field penetration depth, and the electron specific heat of $\text{YBa}_2\text{Cu}_3\text{O}_{7-x}$.

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