

# Gallium-germanium neutrino experiment and conservation of electric charge

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A limit is imposed on the lifetime of  $^{71}\text{Ga}$  with respect to decay involving a violation of the conservation of electric charge. This limit can be established in a gallium-germanium neutrino experiment. The lifetime of the nuclei for which such a decay is possible is shown to be related to a possible deficiency of electrons in the universe.

Feinberg and Goldhaber<sup>1</sup> were the first to point out that if two nuclei  $A_z$  and  $A_{z+1}$  have a mass difference  $M_{\text{nucl}}(A_z) - M_{\text{nucl}}(A_{z+1})$  which is less than the mass of an electron, the nuclide  $A_z$  could be used to test the conservation of electric charge, since the decay of such a nucleus,



where  $X^0$  is an electrically neutral particle or group of particles with a mass smaller than the mass difference of the nuclei  $A_z$  and  $A_{z+1}$  (e.g.,  $\gamma$  or  $\nu + \bar{\nu}$ ), is not forbidden by anything except conservation of electric charge.

One such nuclide is  $^{71}\text{Ga}$ , which is of particular interest in connection with two gallium-germanium experiments which are currently being set up to detect solar neutrinos.<sup>3,4</sup> The results of these experiments will substantially improve the limit on the half-life of the decay of  $^{71}\text{Ga}$  in reaction (1) from the value found in Ref. 5. For example, if the measured effect is the value predicted by the standard solar model<sup>6</sup> without boron neutrinos, which are not observed in a chlorine-argon experiment,<sup>7</sup> this limit would be

$$T_{1/2}(^{71}\text{Ga}) \gtrsim 2 \times 10^{26} \text{ yr.}$$

The nucleus  $A_{z+1}$ , which is formed in reaction (1) decays back into  $A_z$ :



The set of processes (1), (2) should result in a buildup of a number of protons in excess of the number of electrons. Over the lifetime of the universe ( $\sim 2 \times 10^{10}$  yr), this excess would be

$$\delta q = \frac{N_p - N_e}{N_p} \approx 10^{10} \frac{\epsilon(A_z)}{\mu(A_z)} \frac{1}{T_{1/2}(A_z)} \quad (3)$$

(provided that the decay half-life of the nuclide  $A_{z+1}$  is much shorter than the life-

time of the universe). In expression (3),  $N_p$  and  $N_e$  are the densities of protons and electrons in the matter of the universe,  $\epsilon(A_z)$  is the weight concentration of the nuclide  $A_z$ ,  $\mu(A_z)$  is its molar weight in grams, and  $T_{1/2}(A_z)$  is the decay half-life of the nucleus  $A_z$  in reaction (1). If an upper limit  $\overline{\delta q}$  could be found on the deficiency of electrons in the universe in some independent way, it would be possible to establish the following limit on the decay half-life of  $^{71}\text{Ga}$ :

$$T_{1/2}(^{71}\text{Ga}) \gtrsim \frac{0.1}{\overline{\delta q}} \text{ yr} \quad (4)$$

(here and below, we are taking the nuclide abundances from Ref. 8).

If we make certain assumptions, we can establish a limit on  $T_{1/2}(^{71}\text{Ga})$  which is more stringent than (4). Let us assume that the nuclear matrix element for process (1) is written as the product of an ordinary  $\beta$ -decay matrix element and a matrix element which does not conserve electric charge. Let us also assume that a term which violates charge conservation appears in the weak-interaction Hamiltonian in the form  $\alpha J_{\mu}^{(+)} \overline{\nu_e} O^i \nu_e J_{\mu}^{(+)}$  which is the charged current of nucleons.<sup>9</sup> Summing the right side of (3) over all nuclides for which reactions (1) and (2) would be possible, we can find a limitation on  $\alpha^2$  [cf. Eq. (54a) in Ref. 9]. An analysis shows that of the more than 20 such nuclides ( $^{41}\text{K}$ ,  $^{55}\text{Mn}$ ,  $^{71}\text{Ga}$ ,  $^{73}\text{Ge}$ ,  $^{87}\text{Rb}$ ,  $^{93}\text{Nb}$ ,  $^{97}\text{Mo}$ ,  $^{109}\text{Ag}$ ,  $^{125}\text{Te}$ ,  $^{131}\text{Xe}$ ,  $^{139}\text{La}$ ,  $^{145}\text{Nd}$ ,  $^{151}\text{Eu}$ ,  $^{153}\text{Eu}$ ,  $^{157}\text{Gd}$ ,  $^{159}\text{Tb}$ ,  $^{170}\text{Er}$ ,  $^{165}\text{Ho}$ ,  $^{179}\text{Hf}$ ,  $^{181}\text{Ta}$ ,  $^{189}\text{Os}$ ,  $^{193}\text{Ir}$ ,  $^{197}\text{Au}$ ,  $^{201}\text{Hg}$ ,  $^{204}\text{Hg}$ ,  $^{205}\text{Tl}$ ) the value of  $\delta q$  would be dominated by  $^{55}\text{Mn}$ . The  $^{71}\text{Ga}$  contribution would be smaller by about an order of magnitude. The other nuclides would make negligible contributions. As a result, the following limitation is imposed on  $\alpha^2$ :

$$\alpha^2 \lesssim 0.5 \overline{\delta q}.$$

For the decay half-life of  $^{71}\text{Ga}$  we then find

$$T_{1/2}(^{71}\text{Ga}) \gtrsim 4/\overline{\delta q}. \quad (5)$$

It follows from limitation (5) that the background created in a gallium-germanium experiment due to nonconservation of electric charge should not exceed  $\sim 6 \times 10^{27} \overline{\delta q}$  solar neutrino units, corresponding to the formation of  $\sim 10^{26} \overline{\delta q}$   $^{71}\text{Ge}$  atoms per day for a 60-metric-ton detector.

At the moment, the study by Barabanov *et al.*<sup>5</sup> gives us our best limitation on the lifetime of  $^{71}\text{Ga}$  with respect to decay (1). The following upper limit can thus be set on the deficiency of electrons caused by processes (1) and (2):

$$\overline{\delta q} \lesssim 2 \times 10^{-23}.$$

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<sup>3</sup>W. Hampel, "The gallium solar neutrino detector," in: Proceedings of AIP Conference. Solar Neutrino and

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<sup>5</sup>I. R. Barabanov, E. P. Veretenkin, V. N. Gavrin *et al.*, *Pis'ma Zh. Eksp. Teor. Fiz.* **32**, 384 (1980) [*JETP Lett.* **32**, 359 (1980)].

<sup>6</sup>J. N. Bahcall and R. K. Ulrich, Inst. for Advanced Study Preprint IASSNS-AST 87/1, 1987.

<sup>7</sup>J. K. Rowley, B. T. Cleveland, and R. Davis Jr., *Astrophys. J. Lett.* **292**, L79 (1985).

<sup>8</sup>A. G. W. Cameron, in: *Essays in Nuclear Astrophysics* (ed. C. Barnes *et al.*), Cambridge Univ. Press, 1982.

<sup>9</sup>J. N. Bahcall, *Rev. Mod. Phys.* **50**, 881 (1978).

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