

Direct experimental observation of discrete correlated single-electron tunneling

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(Submitted 6 March 1987)

Pis'ma Zh. Eksp. Teor. Fiz. **45**, No. 8, 389–390 (25 April 1987)

Periodic oscillations observed on the current-voltage characteristics of granular tunnel junctions can be described well by a theory of a correlated tunneling of single electrons. This effect can be utilized to develop an ultrasensitive (subelectron) electrometry.

The recent prediction¹ of coherent “single-electron” oscillations with a frequency $f = I/e$, which should be excited by a direct current I in a small-area tunnel junction at a low temperature, has rekindled interest in the effect that underlies these oscillations: a correlated discrete tunneling of single electrons. This effect has previously been observed^{2,3} only indirectly, from its effect on the collective characteristics of systems of a large number of metal granules. In the present letter we report a direct observation of correlated tunneling in a system of two junctions formed by a single metal granule of submicron dimensions.

Figure 1 is a schematic sectional view of the thin-film structures, with an area of $20 \times 20 \mu\text{m}^2$, which we studied. Indium granules with an average diameter on the order of 100 nm are formed between two solid films of lead alloys in such a way that they are separated from the lower film by an In_2O_3 tunnel barrier with a conductivity of about 10^6 S/cm . The granules are separated from the upper film by the same barrier and also by an insulating layer of SiO_2 , about 50 nm thick, which completely covers most of the granules. Consequently, about half of the samples are essentially nonconducting (have a resistance $R \gtrsim 10^7 \Omega$), while the others have a resistance $R \gtrsim 10^5 \Omega$, which corresponds to conduction through a small number ($1-10^2$) of granules. At liquid-helium temperatures such samples exhibit the characteristic suppression of the tunnel current (i.e., a maximum of the derivative $R_d \equiv dV/dI$) in the range $|V| < V_\Sigma$, with a width $2V_\Sigma$ significantly greater than the value $2V_g \cong 2\Delta_{\text{pb}}(T)/e$, associated with the superconductivity of solid films. Furthermore, as the current is increased, the $I(V)$ curves tend toward essentially linear asymptotes, which are shifted precisely $2V_t = 2V_\Sigma - 2V_g$ (Fig. 2a). This effect, which has been seen previously,² can be explained in a natural way¹⁻⁵ in terms of a “blockade” of the tunneling that arises because of electrostatic effects. In accordance with Ref. 2, the $I(V)$ and $R_d(V)$ curves for most of the samples are monotonic.

For a few of the samples, however, these curves exhibit large oscillations ($\Delta R_d/R_d$ up to $\pm 10\%$) with a period ΔV which remains strictly constant for a given sample and with an amplitude which falls off slowly with increasing V (Fig. 2a). At liquid-helium temperatures, this effect is completely reproducible. After the sample is heated to room temperature and then re-cooled, however, the oscillation phase fre-

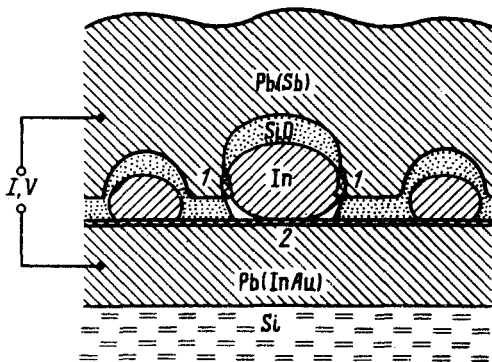


FIG. 1. Sectional diagram of the test structures.

quently shifts with respect to the origin (while the period ΔV is exactly conserved).

This is precisely the behavior that is predicted^{4,5} by the theory of discrete single-electron tunneling in the case in which the conductivity of the system is dominated by tunneling across two junctions (1 and 2 in Fig. 1), which connect electrodes through a single granule. Each oscillation period corresponds to a change in the average charge of the granule by an amount e ; the size of the period in the voltage is $\Delta V = 3/C_1$,

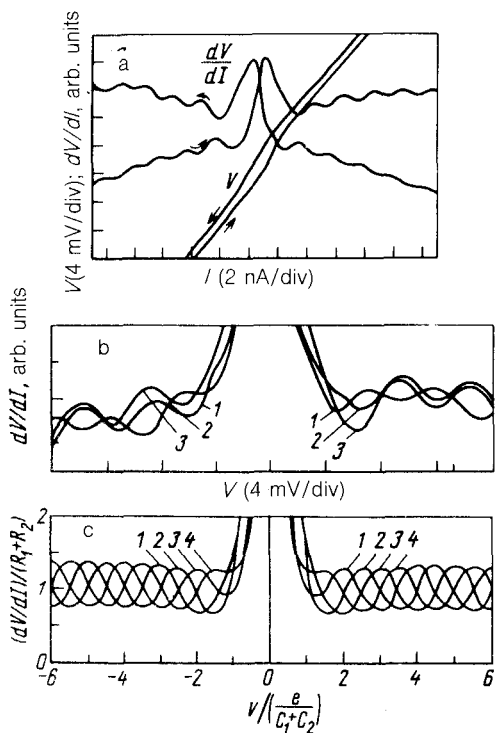


FIG. 2. The voltage V and the derivative dV/dI versus the current I or versus V . a: Experimental (Sample $M 147-V$). The average slope of the $R_d(I)$ curves and their shift upon the reversal of measurement direction are of instrumental origin. b: Three experiments with the same sample at $T = 4.2$ K, alternated with cycles of heating the sample to $T \approx 300$ K. c: Theoretical predictions. 1— $Q_0/e = 0$; 2—0.25; 3—0.5; 4—0.75 ($C_2/C_1 = 1$).

where C_1 is the capacitance of the junction with the lower conductance ($G_1 \ll G_2$). According to the theory, the times at which an electron tunnels across junctions 1 and 2 are mutually correlated, although successive events of tunneling across one of the junctions may be uncorrelated. Figure 2c shows the functional dependence $R_d(V)$ calculated from the theory of Ref. 5 for the experimental parameter values $C_1 = e/\Delta V = 3.17 \times 10^{-7} F$, $T = 4.2 K$, and $\Delta(T) = 1.2 meV$, with the one adjustable parameter C_2/C_1 , for several values of the parameter Q_0 , which is the fractional part of the net electric charge of a granule. A comparison of Figs. 2b and 2c leaves hardly any room for doubting the validity of the interpretation offered above.

Furthermore, our structures can be used in a study of subelectron changes in Q_0 , which apparently involve a movement in the tunnel barriers of isolated impurities with net charges over a distance on the order of interatomic distances.³ Preliminary studies showed that at liquid-helium temperatures Q_0 remains constant within at most $3 \cdot 10^{-2}e$ for tens of minutes during the application of voltages of tens of millivolts. At room temperature, however Q_0 may vary rapidly, by an amount on the order of e , as a result of the application of the voltage of only a fraction of a millivolt. Both of these results are extremely favorable for possible practical applications of this effect.⁵

We wish to thank A. G. Odintsov for assistance in synthesizing the samples.

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Translated by Dave Parsons