

Structural and superconducting phase transitions in V_3Si samples with various defect concentrations

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The temperature of the structural transition, T_M , and that of the superconducting transition, T_c , have been studied experimentally as functions of the defect concentration in the intermetallic compound V_3Si . It is found that T_M falls off monotonically with increasing defect concentration, while T_c goes through a maximum where the temperatures of the structural and superconducting transitions coincide.

In the temperature interval 17–30 K, V_3Si of stoichiometric composition undergoes a structural transition,¹ apparently second-order, from a cubic phase to a tetragonal phase with $c/a = 1.0024$. The physical reasons for the structural transition and its effect on the superconducting properties of V_3Si and other A-15 compounds have been the subject of extensive discussions,^{2–8} primarily from the theoretical standpoint. It has been suggested that the structural transition corresponds to the appearance⁹ or intensification¹⁰ of superconductivity. In this letter we report an experimental study of the structural and superconducting phase transitions in V_3Si samples with various defect concentrations. We detected both transitions from the corresponding jumps in the specific heat.

The composition of these V_3Si samples was nearly stoichiometric. The defect concentration was determined by measuring the residual electrical resistance. The resistivity ratios $\alpha = \rho_{300\text{ K}}/\rho_{18\text{ K}}$ were 50, 31, 21, 19, 16, 14, 12, 11, and 3.4. The specific heat of samples with a mass on the order of 1 g was measured by an adiabatic method¹¹ over the temperature interval 4.5–30 K both without a magnetic field and with a magnetic field of induction up to 20 T. The temperature was measured with a carbon resistance thermometer,¹² with a correction for the effect of the magnetic field.¹³ The operation of the calorimeter was tested by measuring the specific heat of a sample of pure copper; it was found that the error in the measurement of the specific heat does not exceed 1% in magnetic fields up to 8 T or 2% up to 20 T.

A magnetic field suppressed the superconductivity of V_3Si , so that we were able to study its specific heat in the normal (non-superconducting) state down to the lowest temperatures. Figure 1 shows the temperature dependence of the specific heat of several V_3Si samples in the normal state (with the superconductivity being suppressed by a magnetic field). On the curves of the specific heat for samples with low defect concentrations we can clearly see jumps due to a structural transition. These jumps actually occur over a short interval along the temperature scale, apparently because of an inhomogeneity of the samples. The temperature at which the jump occurs, T_M , ranges from 20.5 K at $\alpha = 50$ to 13 K at $\alpha = 12$. The curves of the specific heat of samples with a high defect concentration ($\alpha \leq 11$) do not exhibit the structural features caused by structural transitions. The rapid evolution of the structural feature upon a slight

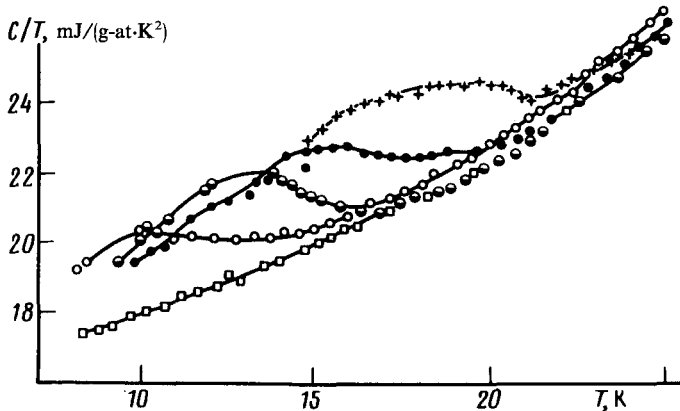


FIG. 1. Specific heat of several V_3Si samples in a strong magnetic field, which suppresses superconductivity. $+ - \alpha = 21$; $\bullet - \alpha = 19$; $\circ - \alpha = 16$; $\bullet - \alpha = 14$; $\square - \alpha = 11$.

change in the defect concentration also shows that these features in the specific heat are in fact due to a phase structural transition. Certain samples exhibit a slight decrease in the transition temperature T_M as the magnetic field is increased. For the sample with $\alpha = 50$ the decrease in T_M in a field of 16 T is 0.9 K, in good agreement with the result of Ref. 14.

According to the present understanding,^{2,3,5} the reason for the structural transition in intermetallic compounds of the A-15 type is a sharp structural feature (a narrow peak, for example) in the electron-state density near the Fermi level. The decrease in the temperature of the structural transition which we observed upon the imposition of a magnetic field and also with increasing defect concentration can be attributed to a

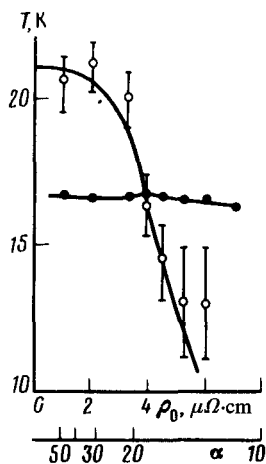


FIG. 2. Temperature of the structural transition, T_M , and temperature of the superconducting transition, T_c , vs the residual resistivity. $\circ - T_M$; $\bullet - T_c$.

spreading of this structural feature in the electron-state density due to either a splitting in the magnetic field or a scattering of electrons by defects.^{15,16}

Figure 2 shows the temperature of the structural transition, T_M , and that of the superconducting transition, T_c , vs the residual resistivity ρ_0 , which is proportional to the defect concentration. The error bars show the width of the structural transition; the width of the superconducting transition corresponds to the size of the circle itself. We see that T_M falls off with increasing ρ_0 , while T_c initially increases, goes through a maximum where the temperatures of the structural and superconducting transitions coincide, and then falls off slowly. The height of this maximum does not exceed 0.15 K, which is less than 1% of T_c . Consequently, at the time of the structural transition there is no significant increase in T_c , and the decrease in T_c with distance from the structural transition is extremely slow.

Whether the structural instability is a primary cause of the high values of T_c in A-15 compounds has been discussed repeatedly in the literature.^{3-5,8} A variety of mechanism have been suggested to explain why T_c increases at or near the structural transition: an intensification of the electron-phonon interaction upon a structural transition,⁹ a softening of the lattice due to a structural instability,⁸ an increase in the electron-state density as the electron spectrum becomes more dielectric,¹⁰ and an interaction of electrons with so-called structural excitations.⁷ The behavior of T_c which we have observed near the structural transition indicates that these mechanisms are apparently not operating in A-15 compounds, and it supports the conclusions of Refs. 5 and 6, according to which the changes in T_c result from changes in the electron-state density accompanying a structural transition.

¹B. W. Batterman and C. S. Barrett, Phys. Rev. Lett. **13**, 390 (1964); Phys. Rev. **145**, 295 (1966).

²Yu. A. Izyumov and É. Z. Kurmaev, Usp. Fiz. Nauk **112**, 193 (1974) [Sov. Phys. Usp. **17**, 44 (1974)].

³L. R. Testardi, M. Weger, and I. Goldberg, Superconducting Compounds with the β -Tungsten Structure (Russ. Transl. Mir, Moscow, 1977).

⁴Yu. A. Izyumov and É. Z. Kurmaev, Usp. Fiz. Nauk. **118**, 53 (1976) [Sov. Phys. Usp. **19**, 26 (1976)].

⁵L. P. Gor'kov and O. N. Dorokhov, Zh. Eksp. Teor. Fiz. **71**, 1934 (1976) [Sov. Phys. JETP **44**, 1014 (1976)].

⁶E. Pytte, Phys. Rev. Lett. **25**, 1176 (1970).

⁷G. M. Vujcic *et al.*, J. Phys. C **14**, 2377 (1981).

⁸L. R. Testardi *et al.*, Phys. Rev. B **3**, 107 (1971).

⁹N. B. Brandt and V. V. Tolmachev, Pis'ma Zh. Eksp. Teor. Fiz. **19**, 439 (1974) [JETP Lett. **19**, 237 (1974)].

¹⁰Yu. V. Kopaev and R. Kh. Timerov, Zh. Eksp. Teor. Fiz. **63**, 290 (1972) [Sov. Phys. JETP **1**, 153 (1972)].

¹¹M. N. Khlopkin, N. A. Chernoplekov, and P. A. Cheremnykh, Preprint IAE-3549/10, I. V. Kurchatov Institute of Atomic Energy, Moscow, 1982.

¹²S. F. Verfolomeev *et al.*, Prib. Tekh. Eksp. No. 1, 262 (1977).

¹³M. N. Khlopkin and P. A. Cheremnykh, Prib. Tekh. Eksp. No. 4, 255 (1981).

¹⁴J. P. Maita and E. Bucher, Phys. Rev. Lett. **29**, 931 (1972).

¹⁵C. J. Williamson, C. S. Ting, and H. K. Fung, Phys. Rev. Lett. **32**, 9 (1974).

¹⁶A. S. Aleksandrov, V. F. Elesin, and M. P. Kazeko, Fiz. Tverd. Tela (Leningrad) **21**, 2062 (1979) [Sov. Phys. Solid State **21**, 1181 (1979)].

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