

## Experimental limitation on the probability for the decay

### $\pi^0 \rightarrow 4\gamma$

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The use of the decay  $K^- \rightarrow \pi^- \pi^0$  as a source of  $\pi^0$  mesons has led to a limitation on the branching ratio for the decay  $\pi^0 \rightarrow 4\gamma$ :  $\text{B.R.}(\pi^0 \rightarrow 4\gamma) < 1.6 \times 10^{-6}$ .

The decay  $\pi^0 \rightarrow 4\gamma$  is allowed by all known selection rules. The branching ratio expected for the decay,  $\text{B.R.}(\pi^0 \rightarrow 4\gamma) = \Gamma(\pi^0 \rightarrow 4\gamma) / \Gamma(\pi^0 \rightarrow 2\gamma)$ , is quite small, and the theoretical predictions are quite uncertain<sup>1,2</sup>:  $\text{B.R.}(\pi^0 \rightarrow 4\gamma) \cong 10^{-9} - 10^{-16}$ . The decay  $\pi^0 \rightarrow 4\gamma$  has also attracted interest, because it is one source of the background in the experiments being carried out to search for the violation of  $C$  invariance in the decay<sup>1</sup>  $\pi^0 \rightarrow 3\gamma$ . A better experimental limitation on the branching ratio for the decay,  $\text{B.R.}(\pi^0 \rightarrow 4\gamma) < 4.4 \times 10^{-6}$ , was found in Ref. 3. In the present letter we report a new value:  $\text{B.R.}(\pi^0 \rightarrow 4\gamma) < 1.6 \times 10^{-6}$ .

The measurements were carried out in parallel with a study for rare modes of in-flight decay of  $\pi^-$  and  $K^-$  mesons at the 70-GeV accelerator of the Institute of High-Energy Physics (the Istra apparatus). We used the decay  $K^- \rightarrow \pi^- \pi^0$  as the source of  $\pi^0$  mesons. The energy of the unseparated beam of  $K^-$  mesons was 25 GeV. The  $\gamma$  rays from the decay of the  $\pi^0$  mesons were detected by a total-absorption Čerenkov spectrometer (Sp), which consists of a matrix of  $24 \times 22$  lead glass counters.<sup>4</sup> The coordinates of the  $K^-$  and  $\pi^-$  mesons were detected by means of hodoscopic photomultipliers and proportional chambers with analog data readout, respectively.<sup>5,6</sup> The

composition and characteristics of the Istra installation are described in more detail in Ref. 4. To single out events with  $n \geq 3$   $\gamma$  rays and to suppress the dominant decay mode  $K^- \rightarrow \pi^- \pi^0 \rightarrow \pi^- 2\gamma$ , we used a system for selecting  $\gamma$ -ray showers in the Čerenkov spectrometer as part of the triggering system.<sup>7</sup>

In selecting candidate events for the decay  $K^- \rightarrow \pi^- \pi^0 \rightarrow \pi^- 4\gamma$  we used the following criteria:

- 1) The presence in the total-absorption Čerenkov spectrometer of four  $\gamma$ -ray showers with a threshold  $E_\gamma > 1$  GeV;
- 2) A total energy evolution of more than 10 GeV in the same Čerenkov spectrometer;
- 3) The presence of only a single track in the hodoscopic photomultipliers and the proportional chambers;
- 4) A minimum distance  $R_{\pi\gamma} > 10$  cm between the projection of the  $\pi^-$  track onto the total-absorption Čerenkov spectrometer and the nearest  $\gamma$ -ray shower.

For all the events that passed these tests, we tested the hypotheses  $K^- \rightarrow \pi^- 4\gamma$  (a 4C fit). We eliminated from further consideration events with  $\chi^2 > 15$ . For the other events, we found the distribution in the effective mass of the four  $\gamma$  rays,  $M_{4\gamma}$ , shown in Fig. 1. The events that lie in the region  $275 \text{ MeV} < M_{4\gamma} < 360 \text{ MeV}$  are the decay  $K^- \rightarrow \pi^- \pi^0 \pi^0$  ( $\tau'$ ); the number of these events after subtraction of the low background is  $N(\tau') = 51\,454 \pm 230$ . This number is used below for a normalization. The arrows in Fig. 1 mark the mass region of the four  $\gamma$ -ray showers which was determined by Monte-Carlo calculations and which corresponds to 95% of the events from the decay  $\pi^0 \rightarrow 4\gamma$ . In this simulation it was assumed that the momenta of the  $\gamma$  rays in the decay  $\pi^0 \rightarrow 4\gamma$  are distributed in phase space. This region contains 38 candidates for the decay  $\pi^0 \rightarrow 4\gamma$ .

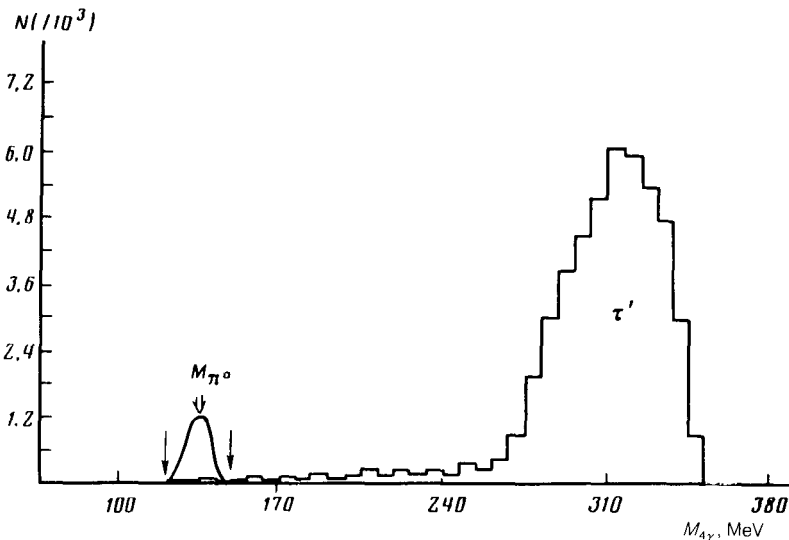


FIG. 1. Distribution in the effective mass of the four  $\gamma$  rays,  $M_{4\gamma}$ . The arrows mark the region of the simulated spectrum of the masses  $M_{4\gamma}$ , which contains 95% of the events from the decay  $\pi^0 \rightarrow 4\gamma$ .

In order to lower the background due to the overload of the total-absorption Čerenkov spectrometer in the central region and the superposition of events that occur at nearly the same time, we then introduced some additional selection rules: a) The lowest of the four energies of the  $\gamma$  rays in the laboratory frame must be  $E_{\gamma\min} > 1.2$  GeV. b) The distance ( $R_{\gamma\min}$ ) from this  $\gamma$  ray to the beam axis must be  $R_{\gamma\min} > 2$  cm. c) We reduced the region of values of the effective mass of the four  $\gamma$  rays to  $129 \text{ MeV} < M_{4\gamma} < 141 \text{ MeV}$ .

These conditions eliminated  $\sim 16\%$  of the events from the decay  $\pi^0 \rightarrow 4\gamma$ . Of the 38 candidate events, none simultaneously satisfied conditions (a), (b), and (c). The upper limit on the number of  $\pi^0 \rightarrow 4\gamma$  decays is thus  $< 2.3$  at a 90% confidence level. The numbers of decays  $\tau'$  and of the decays  $\pi^0 \rightarrow 4\gamma$  for a particular number of decayed  $K^-$  mesons,  $N_k$ , are

$$N(\tau') = N_k \text{ B.R.}(\tau') \epsilon_{\text{geom}} \epsilon_{\chi^2} \epsilon,$$

$$N(\pi^0 \rightarrow 4\gamma) = N_k \text{ B.R.}(K^- \rightarrow \pi^- \pi^0) \text{ B.R.}(\pi^0 \rightarrow 4\gamma) \epsilon'_{\text{geom}} \epsilon'_{\chi^2} \epsilon',$$

where  $\epsilon_{\text{geom}}$  is the geometric efficiency at which the process is detected,  $\epsilon_{\chi^2}$  is the effective cutoff along the  $\chi^2$  scale, and  $\epsilon$  is the resultant efficiency based on the selection criteria and the detection efficiency. After some simple manipulations, we find

$$\text{B.R.}(\pi^0 \rightarrow 4\gamma) = \frac{N(\pi^0 \rightarrow 4\gamma)}{N(\tau')} \frac{\text{B.R.}(\tau')}{\text{B.R.}(K^- \rightarrow \pi^- \pi^0)} \frac{\epsilon_{\text{geom}}}{\epsilon'_{\text{geom}}} \frac{\epsilon_{\chi^2}}{\epsilon'_{\chi^2}} \frac{\epsilon}{\epsilon'}.$$

For the values  $\frac{N(\pi^0 \rightarrow 4\gamma)}{N(\tau')} < 4.5 \times 10^{-5}$ ,  $\frac{\text{B.R.}(\tau')}{\text{B.R.}(K^- \rightarrow \pi^- \pi^0)} = 0.08$ , and (the following were found through the simulation)

$$\frac{\epsilon_{\text{geom}}}{\epsilon'_{\text{geom}}} = 0.37; \quad \frac{\epsilon}{\epsilon'} = 1.2; \quad \frac{\epsilon_{\chi^2}}{\epsilon'_{\chi^2}} = 1,$$

we find  $\text{B.R.}(\pi^0 \rightarrow 4\gamma) < 1.6 \times 10^{-6}$ .

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<sup>1</sup>A. V. Tarasov, *Yad. Fiz* **5**, 445 (1967) [*sic*].

<sup>2</sup>R. L. Schult and B. L. Young, *Phys. Rev. D* **6**, 1988 (1972).

<sup>3</sup>Averbach *et al.*, *Phys. Rev. B* **90**, 317 (1980).

<sup>4</sup>B. H. Bolotov *et al.*, Preprint P-0428, Institute of Nuclear Research, Academy of Sciences of the USSR, Moscow, 1985.

<sup>5</sup>V. G. Vasil'chenko *et al.*, Preprint 78-16, Institute of High-Energy Physics, Serpukhov, 1978.

<sup>6</sup>V. N. Bolotov *et al.*, *Nucl. Instrum. Methods* **A227**, 287 (1985).

<sup>7</sup>S. N. Gninenko *et al.*, Preprint P-0427, Institute of Nuclear Research, Academy of Sciences of the USSR, Moscow, 1985.