

# Strange diffusion of $\mu^+$ mesons in bismuth

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The temperature dependence of the  $\mu^+$ -meson spin relaxation rate  $\Lambda(T)$  in bismuth was measured at  $T < 250$  K. One way to interpret the experimental  $\Lambda(T)$  dependence is to assume coherent quantum diffusion of the  $\mu^+$  meson at  $T < 80$  K.

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The temperature dependence of the  $\mu^+$ -meson spin relaxation rate  $\Lambda_{\text{Cu}}(T)$  in copper was measured in<sup>[1,2]</sup>. This dependence is flat-topped, i. e.,  $\Lambda_{\text{Cu}}(T) = \text{const}$  at  $T < 80$  K and is a monotonically decreasing function ( $\Lambda_{\text{Cu}}(T) \rightarrow 0$ ) at higher temperatures. The rate of relaxation of the  $\mu^+$ -meson spin on the flat top of  $\Lambda_{\text{Cu}}(T)$  is determined by the dipole relaxation of the  $\mu^+$ -meson spin on account of the interaction of the magnetic moment of the  $\mu^+$  meson and the copper nuclei. The drooping part of  $\Lambda_{\text{Cu}}(T)$  at increased temperature can be naturally attributed to diffusion of the  $\mu^+$  meson over the copper lattice.

We have measured the temperature dependence of the  $\mu^+$ -meson spin relaxa-

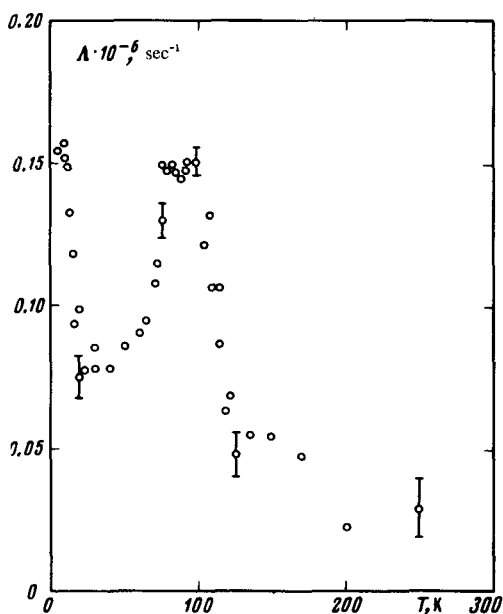


FIG. 1. Temperature dependence of the  $\mu^+$ -meson spin relaxation rate in bismuth.

tion rate  $\Lambda(T)$  in bismuth; this dependence differs substantially from the above-described  $\Lambda(T)$  in copper. The employed bismuth sample was prepared from polycrystalline material with less than 0.01% impurities. The work was performed with the cyclotron of the JINR Laboratory of Nuclear Problems in Dubna. The obtained experimental  $\Lambda(T)$  dependence is shown in Fig. 1. The values of  $\Lambda$  shown in Fig. 1 were determined by the maximum-likelihood method from the damping of the precession amplitude of the  $\mu^+$ -meson spin in a transverse magnetic field  $H=70$  Oe, by comparing the experimental precession curves with the theoretical expression

$$N(t) = N_0 e^{-t/\tau_0} [1 - a e^{-\Lambda^2 t^2} \cos \omega t]. \quad (1)$$

Here  $N(t)$  is the number of positrons of the  $\mu^+ \rightarrow e^+$  decay, emitted in the direction of the primary  $\mu^+$ -meson polarization corresponding to the time  $t=0$ ;  $\tau_0 = 2.2 \times 10^{-6}$  sec is the lifetime of the  $\mu^+$  meson;  $a$  is the experimental asymmetry coefficient of the angular distribution of the  $\mu^+ \rightarrow e^+$  decay electrons;  $\omega = eH/m_\mu c$  is the Larmor precession frequency of the  $\mu^+$ -meson spin in a field  $H=70$  Oe. It follows from (1) that the  $\mu^+$ -meson spin relaxation rate  $\Lambda$  is determined under the assumption that the  $\mu^+$ -meson polarization has a Gaussian time dependence  $P(t) = \exp(-\Lambda^2 t^2)$ . A Gaussian  $P(t)$  should indeed be observed in the absence of  $\mu^+$ -meson diffusion.<sup>[2]</sup> For a diffusing  $\mu^+$  meson, the  $P(t)$  dependence is more complicated, and with increasing diffusion rate (i.e., at high temperatures) it approaches asymptotically an exponential function of the time.<sup>[3]</sup> Figure 1 shows  $\Lambda(T)$  plotted, for the sake of uniformity, under the assumption that  $P(t) = \exp(-\Lambda^2 t^2)$ , i.e., that it agrees with (1) in the entire investigated temperature interval.

It is seen from Fig. 1 that  $\Lambda(T)$  in bismuth is nonmonotonic and its description is not as simple as that of the corresponding relation for copper.<sup>[2]</sup> The appreciable decrease of the relaxation rate at temperatures  $T \gtrsim 150$  K can be naturally explained as being due to the rapid diffusion of the  $\mu^+$  meson at high temperatures. To explain the minimum of the  $\Lambda(T)$  curve at  $T \approx 25$  K, however, it is necessary to take into account some other processes. One of the possible explanations of this effect may be the relatively slow (at low temperatures) diffusion of the  $\mu^+$  meson in the more stable equilibrium positions. At higher temperatures, this diffusion becomes too fast to be observed in the experiment.

Another possible explanation of the minimum of  $\Lambda(T)$  at  $T \sim 25$  K is the coherent quantum diffusion<sup>[4,5]</sup> of the  $\mu^+$  meson in bismuth at  $T < 80$  K. In contrast to ordinary diffusion, the average particle velocity in quantum diffusion increases with decreasing temperature. In quantum diffusion the relaxation rate of the  $\mu^+$ -meson spin should therefore decrease with decreasing temperature, as in indeed observed at  $T < 80$  K. To explain the experimentally observed increase of  $\Lambda(T)$  at  $T < 20$  K it must be assumed that the  $\mu^+$  meson can be captured by a dislocation or by some other bismuth crystal-lattice defect, which becomes possible precisely at low temperatures, when the coherent diffusion velocity increases sharply.

We note in conclusion that the limiting temperature  $T_0$ , below which coherent diffusion in the crystal is possible, depends essentially on the mass  $m$  of the diffusing particle. This dependence is determined by the expression<sup>[4]</sup>  $T_0 \sim \Theta(\Delta/\Theta)^{1/7}$ . Here  $\Theta$  is the Debye temperature of the given crystal;  $\Delta = \exp(-\text{const}\sqrt{m})$  is the width of the coherent band for the diffusing particle of

mass  $m$ . Quantitative estimates show that for a  $\mu^+$  meson in metal we have  $\Delta \sim 1$  K and  $T_0$  does not differ substantially from the Debye temperature  $\Theta$ . Thus, for bismuth, where  $\Theta = 120$  K we have  $T_0 \sim 50$  K, in full agreement with the possibility of observing coherent  $\mu^+$ -meson diffusion in this metal at  $T < 80$  K. Notice should also be taken of the monoisotopic composition of bismuth, which also favors the coherent processes.

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