

Effect of iron impurity on the superconducting transition temperature and magnetic properties of some ternary molybdenum sulfides

N. E. Alekseevskii, G. Vol'f, N. M. Dobrovol'skii, Yu. F. El'tsev,
V. M. Zakosarenko, and V. I. Tsebro

*Institute of Physics Problems, USSR Academy of Sciences
and P. N. Lebedev Physics Institute, USSR Academy of Sciences*

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A strong suppression of superconductivity by an iron impurity and large effective magnetic moments per impurity atom in the ternary molybdenum sulfides such as PbMo_6S_8 , which can cause a strong indirect exchange, indicate that the investigated compounds may be ferromagnetically unstable.

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The ternary molybdenum chalcogenides (TMC) of the formula MMo_6X_8 ($\text{X} = \text{S}, \text{Se}$), where, as is well known, M can be a number of different elements (see, for example, Refs. 1–5), are some of the more interesting and currently intensively studied classes of superconducting components. The effect of ferromagnetic impurities of 3d-elements on the superconducting properties of TMC has been studied very little. There is some evidence that the Fe and Mn impurities strongly suppress superconductivity of SnMo_6S_8 .^[4,6,7]

In this paper we investigate the effect of iron impurity in the region of low concentrations on the superconducting transition temperature T_c and magnetic susceptibility of the ternary sulfides MMo_6S_8 in the normal state, where $\text{M} = \text{Pb}, \text{Sn}, \text{Cu}$, and Ag . The choice of compounds was determined by the fact that when the difference between the critical temperatures T_c is small, the values of $\partial H_{c2}/\partial T$ for the compounds with Pb and Sn are appreciably higher than those for the compounds with Cu and Ag (see Table I). The samples of the $\text{Fe}_x\text{MMo}_6\text{S}_8$ compounds were prepared by a direct synthesis from powders of the starting components and then annealed in the same way as earlier (see, for example, Refs. 4 and 7). The maximum concentration of iron x in the basic formula for all compounds was ≤ 0.05 . The superconducting transition temperature was determined inductively by using ac current. The magnetic susceptibility was measured in the magnetic field of the superconducting solenoid by a string magnetometer whose operating principle was described elsewhere.^[8]

The dependence of superconducting transition temperature on the concentration of the iron impurity measured by us is shown in Fig. 1 for all four systems. The results are given in the T_c/T_{c0} coordinates, where T_{c0} is the superconducting transition temperature of the pure sample. In the investigated region of concentrations it can be seen that T_c for all compounds decreases linearly with increasing concentration of the impurity: the values of $\partial T_c/\partial c$ and $(1/T_c)(\partial T_c/\partial c)$ per at.% of the impurity are given in Table I.

The measurements of the magnetic susceptibility χ of the samples without mag-

TABLE I.

Composition	T_c, K	$\frac{\partial H_{c,2}}{\partial T} \frac{k\bar{G}}{K}$	$\frac{\partial T_c}{\partial c} \frac{K}{\text{at.}\%}$	$\frac{1}{T_c} \frac{\partial T_c}{\partial c}$ (1/at.%)	$\chi (300 K)$ 10^{-6}CGS/g	$\gamma \frac{\text{mJ}}{K^2 \cdot \text{mole}}$	$\gamma_{\Delta c} \frac{\text{mJ}}{K^2 \cdot \text{mol}}$	μ_{eff}/μ_B
$PbMo_6S_8$	13.5	50	28	2.1	0.37	—	101	5
$SnMo_6S_8$	11.4	37	26	2.3	0.46	106	85	4.2
$Cu_{1,8}Mo_6S_8$	10.9	15.5	7.8	0.72	0.26	63	46	3
$AgMo_6S_8$	8.5	7	4.6	0.54	0.09	32	27	2.5

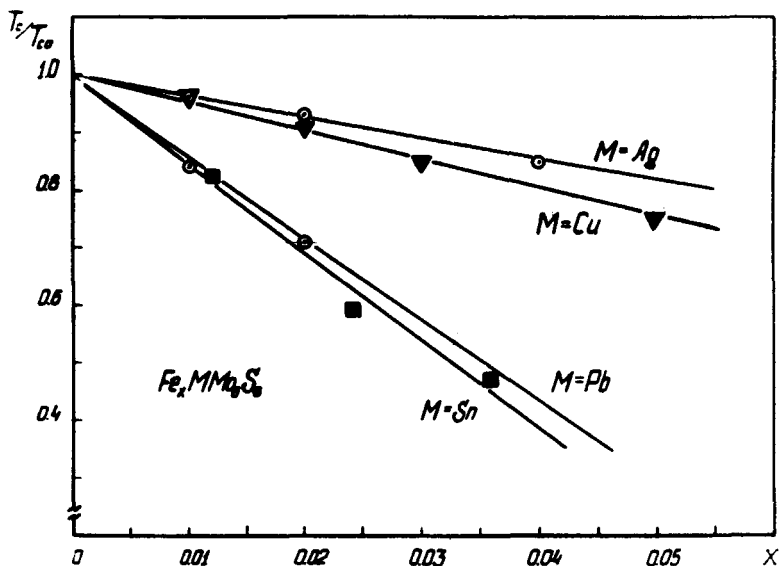


FIG. 1. Dependence of the superconducting transition temperature T_c/T_{c0} on the iron concentration.

netic impurities showed that χ is almost independent of temperature. As the temperature decreased from 300 K to T_c , the magnetic susceptibility, whose values are given in Table I, increased not more than 15% for all the samples. At $\sim T_c$ the magnetic susceptibility of the iron-containing samples increases significantly with increasing concentration of the iron and decreases significantly with increasing temperature, so that the temperature dependence of the reverse susceptibility $\chi^{-1}(T)$ is almost linear. Figure 2 shows the dependence of $\chi^{-1}(T)$ for $Fe_xPbMo_6S_8$ samples with different con-

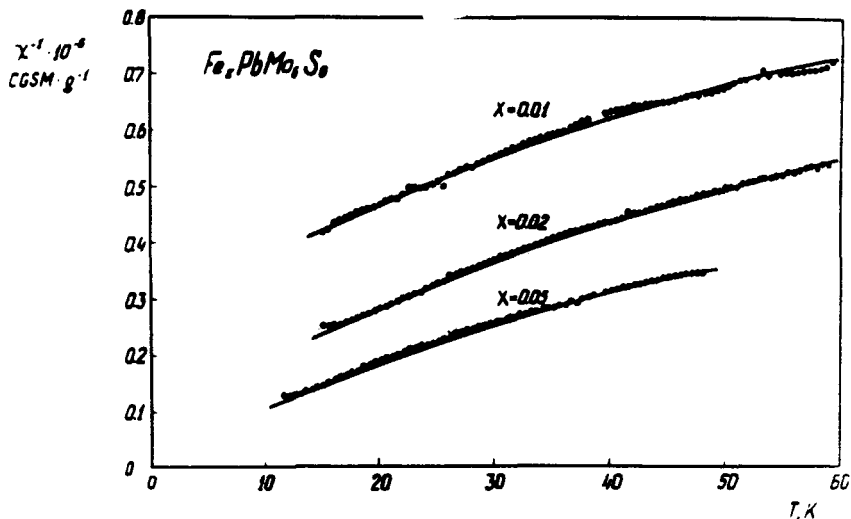


FIG. 2. Temperature dependence of the reverse susceptibility of the $Fe_xPbMo_6S_8$ samples for different concentrations of the iron impurity.

centrations of iron.

The experimental curves $\chi(T)$ for the samples of all the compounds were approximated by an expression of the type

$$\chi(T) = \frac{c}{T - \Theta} + \chi_0, \quad (1)$$

where the parameters c , Θ , and χ_0 were determined by the least-squares method. The solid curves in Fig. 2 correspond to the given approximation. The effective magnetic moments per impurity atom, expressed in Bohr magnetons, which were obtained from the Curie constants according to the formula $\mu_{\text{eff}} = (3k_B c/N)^{1/2}$, are given in Table I. The values of the parameter Θ for all compounds were negative and did not exceed 10 K in absolute value and the values of χ_0 were close to those for the susceptibility of the pure samples.

The superconducting transition temperature generally does not decrease significantly in high-temperature superconducting compounds because of the addition of an iron impurity. Thus, for example, the critical temperature T_c of the intermetallic compounds with A15 structure can be reduced $\sim 10\%$ by introducing up to 5 at.% Fe.¹⁹ In our case, an addition of only 0.5 at.% Fe reduced by 20% the transition temperature of AgMo_6S_8 , $\partial T_c/\partial c$, i.e., the effect is more than an order of magnitude stronger than for superconductors with A15 structure.

It can be seen from Fig. 1 and from the data in Table I that the values of $\partial T_c/\partial c$, which for sulfides with Pb are close to those with Sn, greatly exceed the corresponding values for the compounds with Cu and Ag, for which the values of $\partial H_{c2}/\partial T$ are much smaller than those for the first two compounds. Experimental studies of the specific heat of TMC conducted by us showed that as a result of electron contribution of the specific heat the γ coefficients for PbMo_6S_8 and SnMo_6S_8 are also much higher than those for $\text{Cu}_{1.8}\text{Mo}_6\text{S}_8$ and AgMo_6S_8 (see Table I). The values of γ in Table I were obtained from the analysis of the temperature dependence of the specific heat above T_c (in the same way as in Ref. 7) and from the measurements of the specific heat in the low-temperature region after the destruction of superconductivity by the magnetic field. Moreover, the values of $\gamma_{\Delta c}$, which were obtained in the BCS approximation from the specific heat jump due to the superconducting transition, are given in a separate column. Since the values of $\partial H_{c2}/\partial T$ and γ are determined by the density of states of the electrons at the Fermi surface $N(0)$, we can see that there is a direct correlation between the obtained values of $\partial T_c/\partial c$, μ_{eff} , and $N(0)$.

Although the values of μ_{eff} for PbMo_6S_8 and SnMo_6S_8 are not as large as those for the solutions of Fe in Pd, they are much larger than those for a pure iron. The large values of μ_{eff} for the Pd-Fe alloys, which are attributable to strong indirect exchange, indicate that the investigated systems may be ferromagnetically unstable. This accounts for the very large values of $\partial T_c/\partial c$. As expected, the magnetic instability is stronger in those systems which have a large density of states $N(0)$ (see, for example, Refs. 10 and 11). Note that, according to the data obtained by us earlier,¹⁷⁾ the γ coefficient increases sharply with increasing concentration of the magnetic impurity, which also confirms in a certain sense the existence of magnetic instability.

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