## Suppression of the strength of spin-isospin excitations in $\beta^+$ decay

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A clearly defined resonance of the Gamow-Teller type has been found in the  $\beta^+$  decay of <sup>147m</sup> Dy. The strength function for the  $\beta^+$  decay of this nucleus is calculated. The experimental probability for the  $\beta^+$  decay is only 36% of the calculated value.

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Observations of the Gamow-Teller resonance in the (p,n) reaction for several heavy and intermediate-weight nuclei<sup>1</sup> have revealed that the total strength observed experimentally for the Gamow-Teller excitations is 0.3–0.5 of the theoretical sum-rule prediction. This resonance is known to be a collective state constructed from particle-hole excitations of the  $(pn^{-1})$  type, which are associated with an angular momentum of  $1^+$  and have an isospin projection  $\mu_{\tau}=-1$ . Several theoretical papers (see Ref. 2, for example) have attributed the suppression of Gamow-Teller excitations to an admixture of  $(\Delta$ -isobar)-(nucleon hole) states in nucleon-(nucleon hole) states.

It is clearly worthwhile to experimentally study the behavior of charge-coupled states, i.e., excitations of the  $(p^{-1}n)$  type with an isospin projection  $\mu_z = +1$ . Such states may be excited in the (n, p) reaction and in  $\beta^+$  decay. A study<sup>3</sup> of the strength functions of  $\beta^+$  decay has shown that in nuclei far from the stability band (with a relatively small neutron excess and a large decay energy) the Gamow-Teller excitations are largely concentrated in a narrow energy interval below the decay energy. Studying these nuclei is the easiest way to study a Gamow-Teller resonance with an isospin projection  $\mu_c = +1$ . Calculations on the excitation of this resonance in the  $\beta$  + decay of several thulium nuclei were reported in Ref. 3. These calculations were carried out in the random-phase approximation with a  $G_{\tau\sigma}(\tau\tau)(\sigma\sigma)$  residual interaction. The experimental values of the probability for  $\beta^+$  decay turned out to be about three times lower than the calculated values. In the thulium nuclei, however, the resonances in the strength functions for  $\beta$  + decay lie near the decay energy. The filling of these resonances by  $\beta^+(\epsilon)$  transitions is slight, and several components of the resonance may lie outside the measurement interval. As a result, it was not possible to correctly evaluate the error in the results in Ref. 3, so that the estimate of the suppression of Gamow-Teller excitations found there was strictly qualitative.

In this letter we report a study of the strength function of the  $\beta^+(\epsilon)$  decay of the nucleus <sup>147m</sup> Dy. The measurements were made by a total- $\gamma$ -absorption spectrometer working on line with the IRIS mass separator. The experimental procedure is described in detail in Ref. 4. We determined the decay half-life  $T_{1/2} = 55.7(5)$  s and the decay energy  $Q_{\epsilon} = 7.18(10)$  MeV, and we constructed the strength function for the

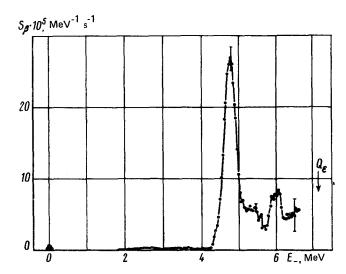


FIG. 1. Strength function of the  $\beta^+(\epsilon)$  decay of  $^{147m}$  Dy.

 $\beta^+(\epsilon)$  decay (Fig. 1). In the strength function we see a narrow peak at 4.84 MeV, flanked on the high-energy side by some fainter peaks. The low position of this resonance (with respect to  $Q_{\epsilon}$ ), the high probability for its occurrence (42% of the decay events), and the rather accurate measurements of  $T_{1/2}$  and  $Q_{\epsilon}$  here make it possible to reliably determine the probability for the excitation of this resonance in  $\beta^+$  decay. This probability corresponds to the value  $\log ft = 3.67(7)$ .

We have also calculated the strength function for the  $\beta$  + decay of this nucleus, using the random-phase approximation. Pairing correlations in the proton system were taken into account in the BCS approximation with a constant  $G_{pair} = 27/$ A MeV. The single-particle energies were taken from Ref. 5. As usual, the spin-isospin interaction constant was taken to be 50/A MeV. From these calculations we found the entire strength of the Gamow-Teller excitations to be concentrated in a single peak. The probability for its excitation corresponds to the value  $\log ft = 3.22(10)$  if the axial vector constant is assumed to have the same value as for the free nucleon (g<sub>4</sub>  $=1.23g_{\nu}$ ). The error in this calculated value is determined by a possible change in the parameters of the average field. Comparison reveals that the excitation probability for Gamow-Teller states with  $\mu_{\tau} = +1$  observed experimentally for the nucleus <sup>147m</sup> Dy is only 0.36(12) of the theoretical value.

Comparison of this result with data from the (p,n) reaction shows that the suppression of the Gamow-Teller excitations in the low-energy region is approximately the same for the isospin projections  $\mu_{\tau}=-1$  and  $\mu_{\tau}=+1$ . It should be noted that a similar suppression has been found<sup>6</sup> for the M1 resonance corresponding to  $\mu_{\tau} = 0$ .

The probability for  $\beta$  decay is related to the matrix element  $M_{GT}$  by

$$\frac{1}{ft} = \left(\frac{G_A}{g_V}\right)^2 |M_{GT}|^2 D^{-1}; \quad D = 6250 \text{ s.}$$

The experimental and theoretical values of ft can be reconciled by assuming that for a complex nucleus the axial vector constant of the weak interaction, G<sub>4</sub>, is not the same as that for the free nucleon,  $G_4$ ; i.e., this constant is renormalized in a nucleus. For the nucleus <sup>147m</sup> Dy in this case we find  $|G_A/g_V| = 0.74(11)$ , so that  $G_A$  is smaller than  $g_A$ by a factor of nearly two.

<sup>1</sup>C. Gaarde et al., in: Fourth International Conference on Nucl. Far from Stability, Helsingor, CERN 81-09, 1981, p. 281.

<sup>2</sup>A. Bohr and B. Mottelson, Phys. Lett. 100B, 10 (1981).

<sup>3</sup>G. D. Alkhazov et al., in: Fourth International Conference on Nucl. Far from Stability, Helsingor, CERN 81-09, 1981, p. 238.

<sup>4</sup>A. A. Bykov et al., Preprint LIYaF-833, Leningrad Institute of Nuclear Physics, Leningrad, 1983. <sup>5</sup>R. R. Chasman, Phys. Rev. C 21, 456 (1980).

<sup>6</sup>G. M. Growley et al., Phys. Rev. C 26, 87 (1982).

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