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It is shown that to obtain agreement with experiment it is necessary to make one subtraction in the dispersion relations for the amplitude of the charge exchange  $\pi^-p \rightarrow \pi^0n$ . The real part of the amplitude and  $d\sigma^{\text{ex}}/dt$  tend in this case to constant values at  $t = 0$  and  $E \rightarrow \infty$ , thus contradicting the model of complex angular momenta.

The differential charge-exchange cross section at  $t = 0$  can be calculated with the aid of the dispersion relations and isotopic invariance if one knows the difference  $\Delta\sigma$  between the total  $\pi^\pm p$  scattering cross sections, namely,

$$d\sigma^{\text{ex}}/dt = 22,5(\Delta\sigma^2 + R^{(-)2}) . \quad (1)$$

Here and below  $d\sigma^{\text{ex}}/dt$  is in  $\mu\text{b}/(\text{GeV}/c)^2$  while  $\Delta\sigma$  and the real part  $R^{(-)}$  of the amplitude are in millibarns.

We have first obtained  $R^{(-)}$  at  $E \gg 1$  GeV from the dispersion relations without subtractions, in the form

$$R^{(-)} = \frac{2}{\pi} \int_0^\infty \frac{dp' p'^2 \Delta\sigma}{E'(E'^2 - E^2)} . \quad (2)$$

Here  $p$  is the momentum and  $E$  the energy of the pion in the laboratory system. For  $\Delta\sigma$  at  $E > 8$  GeV we used the parametrization

$$\Delta\sigma = \sigma_- - \sigma_+ = Q/E^A . \quad (3)$$

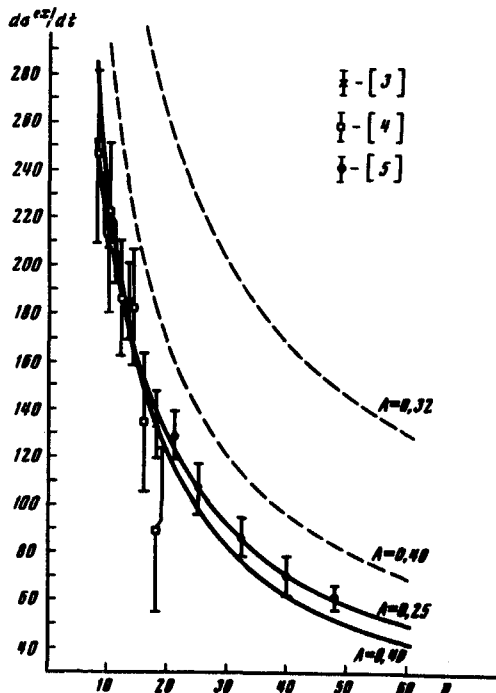


Fig. 1

The calculations were performed for many values of the parameter  $A$  in the interval from 0.25 to 0.40, which covers with large margin the range of variation of this parameter,  $0.32 \pm 0.02$ , obtained in experiments by the Brookhaven group [1] and with the Serpukhov accelerator [2]. The parameter  $Q$  at fixed  $A$  was obtained from the minimum  $\chi^2$  condition.

The correction to formula (3) for  $\Delta\sigma$ , which arises in the theory of complex angular momenta when account is taken of the contributions from the cuts, reduces effectively to a redefinition of the constant  $A$  and does not influence our results.

The values of  $d\sigma^{\text{ex}}/dt$  calculated from formula (1) approach the experimental points [3 - 5] with increasing  $A$ , but even at  $A = 40$  there is no agreement with experiment (Fig. 1, dashed lines).

It is thus obvious that to determine  $R^{(-)}$  it is necessary to use dispersion relations with subtraction

$$R^{(-)} = G + \frac{2E^2}{\pi} \int_0^{\infty} \frac{dp' \Delta\sigma}{E'(E'^2 - E^2)} \quad (4)$$

By suitably choosing the subtraction constant  $G$  we can obtain good agreement between the calculated  $d\sigma^{\text{ex}}/dt$  and the experimental data for all values of  $A$  in the interval from 0.25 to 0.40 (Fig. 1, solid lines).

The discrepancy between (2) and experiment can be attributed to the fact that actually  $\Delta\sigma \rightarrow \text{const}$  as  $E \rightarrow \infty$ . Such a possibility, which corresponds to violation of the Pomeranchuk theorem, is considered in [6].

On the other hand, if  $\Delta\sigma$  does satisfy the Pomeranchuk theorem and is described by relation (3), then it follows from our calculations that  $R^{(-)} \rightarrow C \neq 0$  as  $E \rightarrow \infty$ . This result contradicts the predictions of the complex angular momentum model.

The dispersion relation (4) reduces in this case to relation (2), in which the constant  $C$  is added on the right.  $C = -1.2$  mb for  $A = 0.32$ . Then  $d\sigma^{\text{ex}}/dt(\infty) = 37 \mu\text{b}/(\text{GeV}/c)^2$ . (At  $A = 0.4$  these quantities are equal to  $-0.5$  mb and  $6.5 \mu\text{b}/(\text{GeV}/c)^2$ , respectively.)

At energies  $10 \text{ GeV} \leq E \leq 60 \text{ GeV}$  the experimental points fit well the curve  $d\sigma^{\text{ex}}/dt = 1590/E^{0.84}$  [5]. This curve is shown dashed in Fig. 2, together with the dispersion curve obtained from (4) at  $A = 0.32$ . In the investigated energy range  $E < 60 \text{ GeV}$ , these curves practically coincide. In the region  $100 \text{ GeV} < E < 220 \text{ GeV}$  the discrepancy between them becomes appreciable, on the order of 25%. The curves subsequently intersect and diverge above 600 GeV.

It is clear from the foregoing calculations that measurement of  $\Delta\sigma$  as well as of the charge-exchange cross section at energies above 60 GeV is extremely important for models of high-energy scattering.

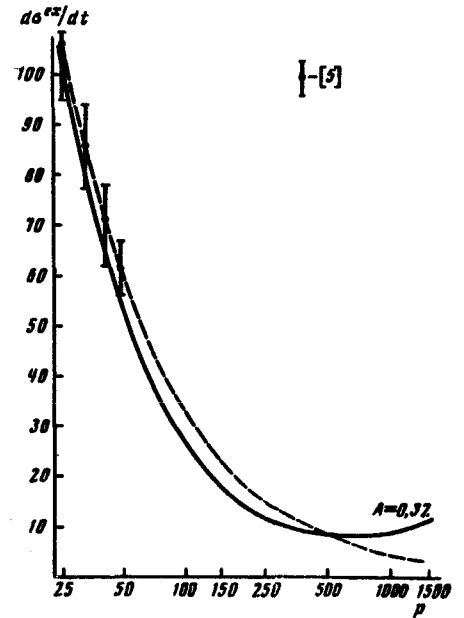


Fig. 2.

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