EXPERIMENTAL CONFIRMATION OF THE INFLUENCE OF LOCALIZED SPIN FLUCTUATIONS (LSF) ON SUPERCONDUCTIVITY

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The theory of Benneman and Garland [1] deals with the influence of the localized spin fluctuations (LSF) on the temperature T of the superconducting transition. It is shown that  ${\tt T}$  of a number of transition-metal alloys is lowered by the LSF by an amount  $^{\rm c}$ 

$$\Delta T_c^{sf} = T_c^{bs} - T_c^{sf},$$

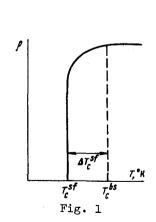
where  $T_c^{bs}$  is the calculated value of  $T_c$  corresponding to allowance for only the band structure of the alloy.  $T_c^{sf} \equiv T_c(expt)$  is the critical temperature with allowance for the LSF.

On the basis of an experimental study of the temperature dependence of the resistivity  $\rho(T)$  in the alloy system  $Ti_x - V_{l-x}$ , we have established a relation between the anomalous properties of these alloys.

On the one hand, the  $\rho(T)$  curves have negative slopes in the temperature range 20 - 300°K. According to many theoretical and experimental studies [2], this is a symptom of LSF.

On the other hand, the transition from the superconducting to the normal state exhibits a typical broadening [3], and traces of superconductivity are preserved up to  $T \simeq 2T_c$ .

Using the scheme of Fig. 1, we have separated a temperature interval  $\Delta T_c$ intermediate between the fully conducting and fully normal state, and set it in



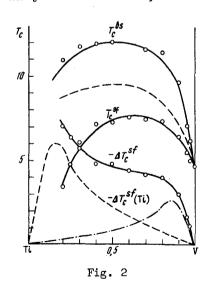


Fig. 1. Scheme of transition into the superconducting state.

Fig. 2. Concentration dependences of  $T_c^{bs}$ , Dashed curves – theoretical plots of  $T_c^{bs}$ 

correspondence with  $\Delta T_c^{sf}$ .

Our experimental results, which are shown in Fig. 2, are in good qualitative agreement with the theory (see Fig. 2 of [1]).

We believe the causes of the quantitative discrepancies between the experimental and theoretical  $\mathbf{T}_{c}^{bs}$ and the noticeable quantitative disparity bebetween the theoretical and experimental  $T_c^{sf}$  to be the following: (1) The experimental values of  $\Delta T_c^{sf}$  contain, in addition to the contribution from the LSF, also a contribution due to thermodynamic fluctuations of the superconducting ordering parameter, which may not be small in this case. (2) In the Benneman-Garland theory, the entire effect of the suppression of  $T_{\rm c}$  is ascribed to LSF on V atoms, and this effect is linear in concentration up to ~10 at.% V.

It is seen from Fig. 2 that a similar situation obtains also for Ti atoms at x < 0.1.

In other words, the experimental  $\Delta T_c^{sf}$  curves contain, besides the contribution  $\Delta T_c^{sf}(V)$ , also a contribution  $\Delta T_c^{sf}(Ti)$ , which is shown schematically by the dash-dot line in Fig. 2. This shows once more that the anomalies described above are due to LSF.

It is interesting to note that the theoretical estimates  $\overset{\text{Tbs}}{\text{c}}_{\text{max}} = 12^{\circ}\text{K}$  in Ti-V alloys, made back in 1961 [4], are in full agreement with our results.

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CROSS SECTION RATIO OF THE REACTIONS  $He^4(\gamma, p)H^3$  AND  $He^4(\gamma, n)He^3$  IN THE RE-GION OF GIANT RESONANCE, AND CHARGE SYMMETRY

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A number of recent publications, by Berman, Fultz, and Kelly [1] and by Berman, Firk, and Wu [2] have been devoted to the reaction  ${\rm He}^4(\gamma,\,n){\rm He}^3$  in the γ-quantum energy range from threshold to 32 MeV. Berman et al [1, 2], using their own data on the total cross section of this reaction and the cross sections given by others [3-5] for the  $(\gamma,\,p)$  reactions, have calculated their ratio as a function of the  $\gamma$ -quantum energy. The average ratio for giant resonance turned out to be 2.0.

Berman et al. used the expressions of Barker and Mann [6], which connect the ratios of the cross sections of these reactions with the amplitudes characterizing the states with isospin T = 0 and T = 1

$$\frac{\sigma_p}{\sigma_n} = \frac{P_p(E_p)}{P_n(E_n)} \left| \frac{\alpha_1 + \alpha_o}{\alpha_1 - \alpha_o} \right|^2$$