

A-dependence of the inclusive distributions of the J/ψ particles

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We measured the differential cross sections $d\sigma/dx$ and $d\sigma/dp_t^2$ of the inclusive production of J/ψ particles by π^- mesons with momentum 43 GeV/c from the nuclei Be, Cu, and W. The dependence of the cross sections on the atomic weight A of the nucleus was approximated by the function A^α . The exponent α increases with increasing p_t^2 and is independent of x within the limits of statistical errors.

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We have previously reported^[1] measurements of the A dependence of the cross section for the production of J/ψ particles by π^- mesons with momentum 43 GeV/c. In the present article we present the distributions of the J/ψ particles with respect to

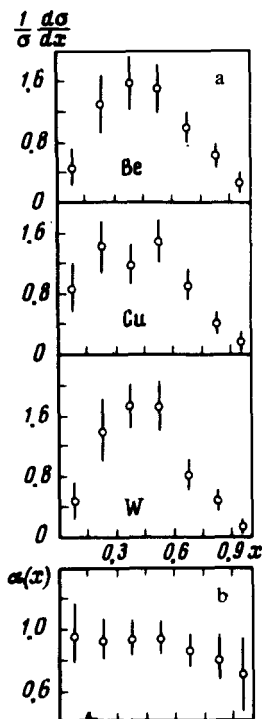


FIG. 1. a) Normalized inclusive cross sections $(1/\sigma)(d\sigma/dx)$ for the production of J/ψ particles on the nuclei Be, Cu, and W; b) dependence of the exponent α on x for the J/ψ particles.

the variables $x = pL^*/p_{max}^*$ and p_t^2 , obtained from the nuclei Be, Cu, and W in the same experiment. The experimental setup (the "Sigma" spectrometer) and the reduction procedures are described in^[1,2].

The reaction $\pi A \rightarrow J/\psi + \dots$ was identified by the $J/\psi \rightarrow \mu^+\mu^-$ decay. To construct the inclusive distributions of the J/ψ particles we selected events with effective pair mass $M_{\mu\mu}$ in the band $|M_{\mu\mu} - M_{J/\psi}| \leq \Gamma_{J/\psi}$, where $M_{J/\psi}$ and $\Gamma_{J/\psi}$ are the measured^[1] mass and width of the J/ψ peak in the effective-mass spectrum of the $\mu^+\mu^-$ pairs. The background under the J/ψ peak was 5–20% in different targets.^[1] It was assumed that this background has the same dependence on x and p_t^2 as the cross section for the production of J/ψ particles.

Figure 1a shows the normalized cross sections $(1/\sigma)(d\sigma/dx)$ for the production of J/ψ particles from different nuclei. Their shapes are similar: the broad maximum at $x \approx 0.4$, observed for copper in^[3,4], is present also in the spectra obtained for beryllium and tungsten. The results of the approximation of the different cross sections $d\sigma/dx$ by the power-law function Ax^α is shown in Fig. 1b. The parameter α is independent of x , within the limits of the statistical errors, and agrees with the value $\alpha = 0.92 \pm 0.06$ obtained by us earlier^[1] for the total inclusive cross sections of J/ψ particle production. Independence of α of x was observed also in the production of J/ψ particles in nA interactions at 300 GeV/c.^[5]

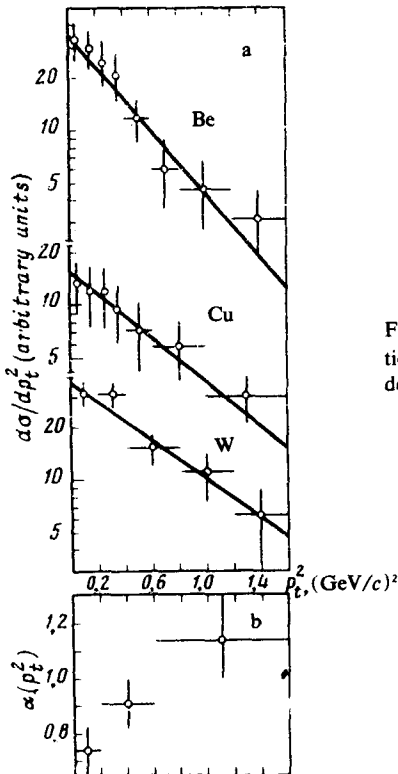


FIG. 2. a) Inclusive cross sections $d\sigma/dp_t^2$ for the production of J/ψ particles from the nuclei Be, Cu, and W; b) dependence the exponent α on p_t^2 for the J/ψ particles.

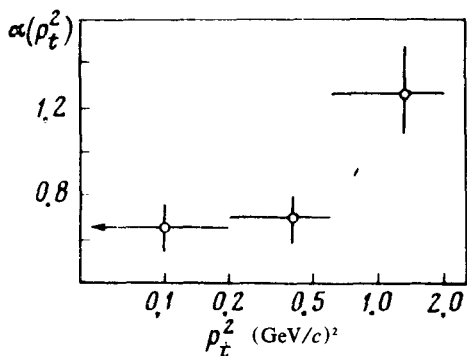


FIG. 3. The same as in Fig. 2b, but with a different subtraction of the background (see the text).

Figure 2a shows the cross sections $d\sigma/dp_t^2$ for the production of J/ψ particles from different nuclei. The straight lines are the results of fitting to the exponential $\exp(-bp_t^2)$. The values obtained for the slope parameter b are $b_{\text{Cu}} = 1.45 \pm 0.25$; $b_{\text{W}} = 1.3 \pm 0.3$ (GeV/c)⁻². The value of b_{Be} agrees within the limits of errors with the one obtained by us earlier,^[2] while b_{Cu} agrees with the value obtained in^[3,4]. Attention is called to the fact that with increasing atomic number of the target the slope b decreases. Figure 2b shows a plot of $\alpha(p_t^2)$ obtained by approximating the cross section $d\sigma/dp_t^2$ by the power-law function A^α . It is seen that α increases from 0.74 ± 0.09 at $p_t^2 = 0.1$ to 1.14 ± 0.14 at $p_t^2 = 1.1$ (GeV/c)². Since the dependence of the cross section for the production of J/ψ particles on p_t^2 can differ from the dependence of the cross section of the production of the $\mu^+\mu^-$ pairs with the same effective mass, we have carried out the following additional analysis. All the events of $\mu^+\mu^-$ pair production were divided into three intervals with respect to p_t^2 and the cross section for the production of the J/ψ particles was determined for each of the intervals by fitting the $M_{\mu\mu}$ mass spectrum in the manner described in^[1]. The resultant function $\alpha(p_t^2)$ is shown in Fig. 3 and is similar to that obtained in Fig. 2b. The exponent α determined from the mass band 2500–2800 MeV adjacent to J/ψ amounts to ~ 0.7 and is independent of p_t^2 within the limits of errors. Thus, the observed increase of α with increase of p_t^2 is due to the production of the J/ψ particle.

A similar behavior of $\alpha(p_t^2)$ was observed in the reactions for the inclusive production π^\pm, K^\pm, p and \bar{p} at 28,^[6] 50–275,^[7] 300,^[8] and 200–400^[9] GeV/c and of K^0 at 10 GeV/c.^[10] It should be noted that no variation of α with p_t^2 was observed in the reaction $nA \rightarrow J/\psi + \dots$ at 300 GeV/c.^[5]

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