

increases at the same time.

The magnetic moment of the first excited state of  $\text{Sn}^{119}$ , calculated from the obtained nuclear gamma-resonance spectra, is  $0.67 \pm 0.01$  nuc. magnetons, in agreement with the data of [6].

The authors are sincerely grateful for very useful discussion to Yu. M. Kagan, who previously called attention to the importance of experiments on the polarization of nuclei in a ferroelectric matrix [7]. It is also a pleasure to thank E. F. Makarov for help with the work, S. S. Kurochkin for the use of the 2048-channel analyzer with which the spectra could be obtained with high resolution, and E. L. Frankevich for help with measuring the conductivity of the samples.

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#### ISOTOPIC EFFECT IN THE FERROELECTRIC $\text{NaH}_3(\text{SeO}_3)_2$

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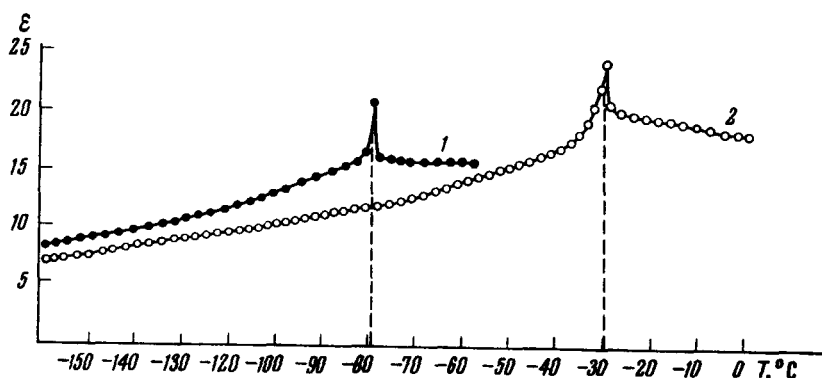
Original submitted 18 February 1965.

The ferroelectric properties of sodium hydroselenite  $\text{NaH}_3(\text{SeO}_3)_2$  (Curie point  $T_C = -79^\circ\text{C}$ ) were investigated by Pepinsky and Vadam [1, 2]. Of great interest in explaining the nature of the spontaneous polarization in this compound is the isotopic effect when the hydrogen is replaced with deuterium.

We report here the results of a study of the temperature dependence of the dielectric constant of powdered samples of  $\text{NaD}_3(\text{SeO}_3)_2$ .

The sodium deuterioselenite was obtained by crystallization from a solution in  $D_2O$  (99.5% pure), in which the calculated amounts of  $Na_2SeO_3$  and  $D_2SeO_3$  were dissolved. The selenious acid  $D_2SeO_3$  was obtained in turn from the reaction between selenious anhydride  $SeO_2$  and heavy water.

Tablets 1 - 2 mm thick and 14 mm in diameter were prepared from the  $NaD_3(SeO_3)_2$  and placed between the electrodes of a parallel-plate capacitor. The capacitance was measured with the aid of a conventional bridge circuit at about 200 kcs. The measurements were made in the temperature range from  $-170^\circ$  to  $0^\circ C$ .



Temperature dependences of the dielectric constants of polycrystalline  $NaH_3(SeO_3)_2$  (1) and  $NaD_3(SeO_3)_2$  (2).

The figure shows the obtained temperature dependence of the dielectric constants of  $NaD_3(SeO_3)_2$ . For comparison, the temperature dependence we obtained for powdered  $NaH_3(SeO_3)_2$  is shown in the same figure. Repeated measurements of tablets prepared in various manners yielded identical results.

The figure shows that the peak of the dielectric constant of polycrystalline  $NaH_3(SeO_3)_2$  lies at  $-79^\circ C$ , which coincides with the previously determined Curie point of this ferroelectric [1, 2]. The temperature dependence of the dielectric constant of  $NaD_3(SeO_3)_2$  duplicates that of  $NaH_3(SeO_3)_2$ , but is shifted  $50^\circ$  in temperature, with a peak at  $-29^\circ C$ , which can be assumed to be the Curie point of the deuterioselenite.

Thus, the shift of the Curie point of the hydroselenite when the hydrogen is replaced by deuterium is of the same order of magnitude as that in  $\text{KH}_2\text{PO}_4$  ( $90^\circ$ ),  $\text{KH}_2\text{AsO}_4$  ( $66.4^\circ$ ), and  $\text{Ag}_2\text{H}_3\text{IO}_6$  ( $40^\circ$ )<sup>[3]</sup>. It can be assumed on this basis that these ferroelectrics have a similar spontaneous-polarization mechanism in which an important role is played by the ordering of the hydrogen bonds.

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#### 75-MICRON LASER

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Original submitted 19 February 1965.

The first laser design was proposed in 1958<sup>[1]</sup> in the form of an ammonia maser operating in the far infrared (several hundred microns). The actually constructed laser, however, operated in the visible region<sup>[2]</sup>.

A proposal was made in 1959 to use a discharge-excited gas as the active medium<sup>[3]</sup>. This method makes it possible to construct lasers operating in a broad range of wavelengths. It is therefore of interest to construct a laser operating in the far infrared and using a gas discharge. In the design of gas lasers for the far infrared attention is usually paid to close levels of the higher states<sup>[4]</sup>. It must be noted that in Xe overlap of the p and d series takes place even for the lower states, so that lasing can be produced with a series of  $3d \rightarrow 2p$  transitions at relatively weak excitation. An attempt to obtain laser action with transitions of the  $2p \rightarrow 3d$  type may fail because of the rapid depletion of the 2p levels resulting from the short-wave spontaneous  $2p \rightarrow 1s$  transitions<sup>[5]</sup>.

The transition with longest wavelength between the states 2p and 3d in Xe is  $2p_5 \rightarrow 3d_5$  (75.5778  $\mu$ ), the states of which have several distinctive features that make lasing possible in this transition. Moreover, we can