

INTERACTION OF NEGATIVE PIONS IN NUCLEAR EMULSION AT ENERGIES 7.5, 17, AND 60 GeV

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 Submitted 16 August 1971  
 ZhETF Pis. Red. 14, No. 7, 405 - 408 (5 October 1971)

NIKFI-R, Ilford-G-5, and NIIKhIMFOTO-BR-2 nuclear emulsions were irradiated in the accelerators at Dubna, Geneva, and Serpukhov by  $\pi^-$  mesons of energy 7.5, 17, and 60 GeV. The characteristics of these emulsions, the scanning technique, etc. are described in [1 - 3]. Scanning along the track yielded 1266, 4234, and 1986 inelastic interactions of pions in the emulsions at the indicated respective energies. The most general characteristics of these interactions are shown in Figs. 1 and 2 and in the table.

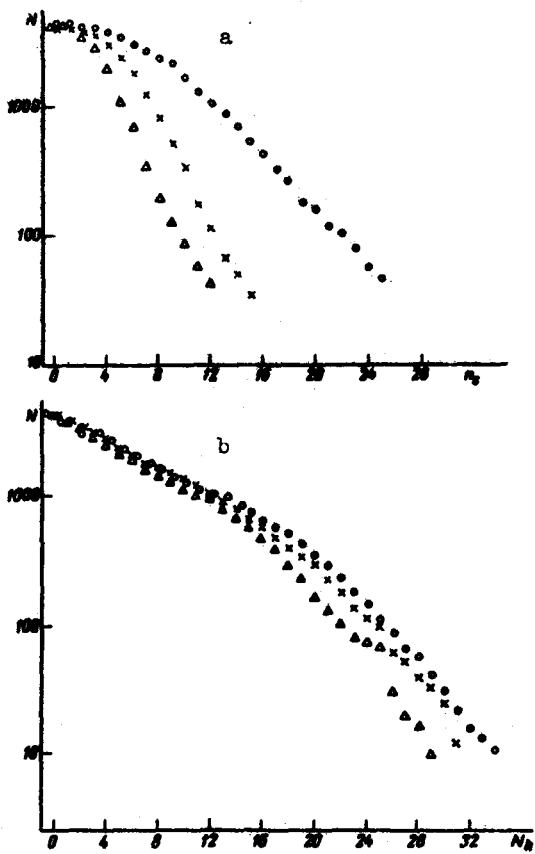


Fig. 1

Figure 1 shows the integral distributions with respect to the multiplicity  $n_s$  and the number of cascade-evaporation particles  $N_h$ ; Fig. 2 shows the dependence of the mean values  $n_s$  and

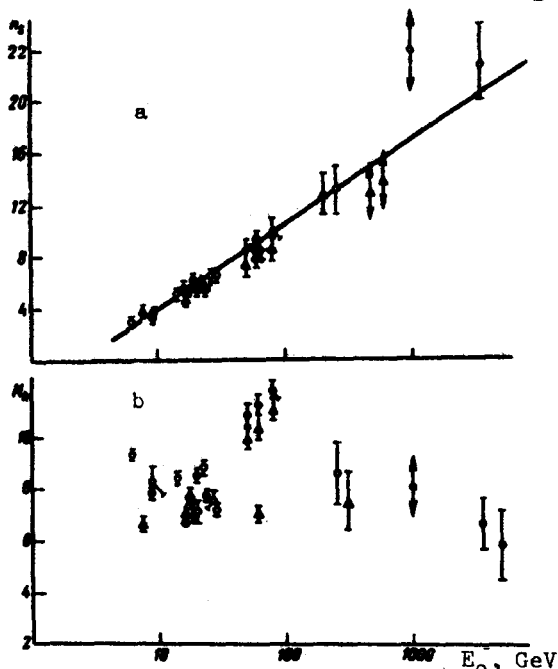


Fig. 2

Fig. 1. Integral distributions with respect to  $n_s$  and  $N_h$ . The points  $\Delta$ ,  $\times$ , and  $o$  correspond to pion-nuclear interactions at energies 7.5, 17, and 60 GeV.

Fig. 2. Dependence of  $n_s$  and  $N_h$  on  $E_0$ .  $\Delta$  and  $\blacktriangle$  - pion-nuclear interactions,  $o$  and  $\bullet$  - proton-nuclear interactions,  $\Delta$  - data of present paper and of [4 - 10],  $o$  - data of [11 - 21],  $\blacktriangle$  and  $\bullet$  - calculated data [22].

$N_h$  on the energy of the primary particles  $E_0$  in accordance with our data and the data of [4 - 10]. Figure 2 shows for comparison analogous plots for proton-nuclear interactions [11 - 21] and the results of calculations by the cascade model with allowance for many-particle interactions [22]. In the region of accelerator energies, we used only the data obtained by analysis of more than 350 nuclear interactions, found by scanning along the track.

The dependence of  $n_s$  on  $E_0$  for pion-nuclear ( $\pi$ -n) and proton-nuclear (p-n) interactions (Fig. 2a) turned out to be practically the same ( $n_s \sim \log E_0$ ), and the dependence of  $N_h$  on  $E_0$  (Fig. 2b) turned out to be different and greatly differing from the results of the calculation [22]. For p-n interactions, there is a tendency towards decreasing  $N_h$  with increasing  $E_0$ , and for  $\pi$ -n interactions there is a tendency towards a very weak growth. It is important to note that the experimental points (Fig. 2b) for p-n interactions lie as a rule higher than for  $\pi$ -n interactions.

The dependences of  $N_h$  and  $n_s$  on  $E_0$  for  $\pi$ -n and p-n interactions can be understood by assuming that the probability of a second collision is smaller for a pion than for a proton, and the multiplicity in pion-nucleon interactions is higher than in proton-nucleon interactions, as is confirmed to some degree by [3, 23, 24]. The probability of a second collision depends on the density of the nucleons in the nucleus and on the particle interaction cross section. Consequently, the difference in  $N_h$  between  $\pi$ -n and p-n interactions, for the same nuclei, is determined by the different cross sections  $\sigma_{\pi N}$  and  $\sigma_{pN}$  of the pion-nucleon and proton-nucleon interactions.

The insignificant dependence of  $N_h$  on  $E_0$  is due to the weak participation in the interaction inside the nucleus by the secondary particles produced in the first collision of the incident particle with the nucleon of the nucleus, where the increase of  $n_s$  with increasing  $E_0$  is due to the development in the nucleus of a cascade of successive interactions of the leading ("conserved") primary particle.

The number of collisions ( $n_{\pi n}$  or  $n_{pn}$ ) can be estimated from the ratio  $n_{pn}/n_{\pi n} = \sigma_{pN}/\sigma_{\pi N}$ , assuming that the difference in the number of collisions of the primary pion and the proton in the nucleus is equal to  $2\Delta N_h$ . In the energy region 9 - 28 GeV, where  $\Delta N_h = 1.1 \pm 0.02$ , and where numerous accelerator data give  $\sigma_{pN} = 1.6\sigma_{\pi N}$ , we get  $n_{\pi n} \approx 3.7$ . Knowing  $n_{\pi n}$  and  $\Delta N_h$ , we can find the ratio of the cross section  $\sigma_{pN}/\sigma_{\pi N} \approx (1 + 2\Delta N_h/n_{\pi n})$ .

Dependence of mean values of  $n_s$  on  $N_h$

| $N_h$ | 7.5 GeV     |              | 17 GeV      |              | 60 GeV       |              |
|-------|-------------|--------------|-------------|--------------|--------------|--------------|
|       | $n_s$       | $N_h$        | $n_s$       | $N_h$        | $n_s$        | $N_h$        |
| 0-1   | 3.30 ± 0.11 | 0.49 ± 0.04  | 3.75 ± 0.06 | 0.45 ± 0.03  | 6.18 ± 0.11  | 0.42 ± 0.03  |
| 2-7   | 3.47 ± 0.08 | 3.82 ± 0.06  | 4.97 ± 0.06 | 3.95 ± 0.05  | 7.61 ± 0.11  | 4.03 ± 0.03  |
| ≥ 8   | 4.08 ± 0.10 | 13.90 ± 0.20 | 6.20 ± 0.10 | 14.44 ± 0.10 | 11.77 ± 0.13 | 15.22 ± 0.15 |
| ≥ 0   | 3.64 ± 0.06 | 6.59 ± 0.07  | 5.24 ± 0.04 | 7.27 ± 0.04  | 9.23 ± 0.07  | 7.02 ± 0.05  |

If the excitation of the nucleus depends on the nature of the particle passing through it, then  $\Delta N_h$  will be determined not only by the cross section of the interaction of the particles in the nucleus, but also by its excitation energy.

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#### MAKING A MEDIUM TRANSLUCENT TO RADIATION EXCITED BY A DIRECTED STREAM OF ATOMIC PARTICLES

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 Submitted 23 August 1971  
 ZhETF Pis. Red. 14, No. 7, 408 - 411 (5 October 1971)

Excitation by atomic collision, unlike excitation by electron impact or photoexcitation, is frequently accompanied by transfer of considerable momentum