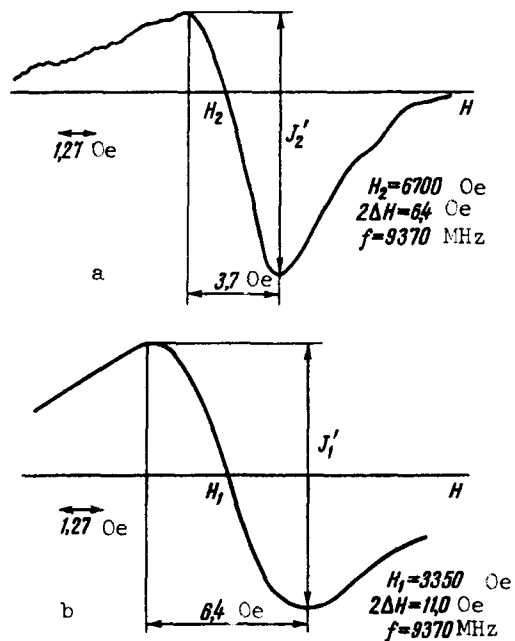


absorption line shape should coincide with the linear FMR line shape.



a - derivative of NFM absorption line, b - derivative of linear FMR absorption line. Intensity ratio  $I_1/I_2 = 10^6$ .

The experimental observation of the NFM was made with a standard radio spectroscopy (RE 1301) at 9370 MHz. The use of a reflex resonator with  $Q = 500$  and with a radiated power on the order of 15 W has made it possible to register, at the sensitivity limit of the spectrometer, the NFM absorption line in a single-crystal sphere made of yttrium iron garnet. The sphere diameter was 3.3 mm. The derivative of the absorption line was recorded on the chart of the automatic plotter (Fig. 1a). The NFM plotting was followed, under the same conditions (orientation and location of the ferrite sphere in the cavity), by a plotting of the derivative of the absorption line for the linear FMR (Fig. 1b). In this case, the high-frequency energy source was the klystron oscillator of the spectrometer, to prevent heating of the sample.

The NFM and linear FMR absorption line widths are 6.4 and 11.0 Oe, respectively. The large experimentally obtained NFM line width is due to distortion of the hf field in the cavity as a result of the strong

microwave absorption in the ferrite.

We see therefore that NFM can be used to study the FMR line shape of large-size samples. In addition, this phenomenon is of considerable interest from the point of view of the study of nonlinear interactions occurring in solids.

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#### LASER WITH FREQUENCY SCANNING DURING THE GENERATION PROCESS

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 Submitted 17 February 1967  
*ZhETF Pis'ma* 5, No. 10, 355-356 (15 May, 1967).

We have observed the effect of frequency variation and the emission kinetics of a laser with a dispersive resonator, in which the maximum  $Q$  was displaced along the frequency scale during the course of the generation, within the limits of the luminescence spectrum of the active medium. For resonators with dispersive prisms [1,2], a controlled frequency variation (scanning) can be accomplished by rotation or angular vibration of the end reflector (Fig. 1), and for resonators with selector-interferometers operating in the transmission mode [3] it can be accomplished by rotating the selector. In the latter case repeated scanning with variable

emission.

Thus, generation with variable emission frequency possesses a number of new useful properties. In fact, tuning of the dispersive resonator during the course of the laser operation leads to an ordering of the temporal and mode character of the generation. The possibility of controlling the kinetics and the spectrum of the generation makes it possible to investigate the spectral structures of the homogeneously and inhomogeneously broadened luminescence bands of condensed media activated with rare-earth ions. We note that for media having spectral inhomogeneities frequency scanning during the generation should increase the energy efficiency as a result of the sequential operation of active centers that do not participate in the generation in the usual "single-frequency" mode.

It is of fundamental interest to investigate the formation of the generation field in the resonator and the genesis of the mode composition of the radiation at those scanning speeds at which the time of resonator tuning, in a frequency interval corresponding to the distance between neighboring lines in the generation spectrum, becomes comparable with or smaller than the growth time of the stimulated-emission avalanche.

Lasers with frequency scanning during the course of generation can be used to investigate the dispersion relations between different effects of interaction between radiation and matter during one generation flash.

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\* A decrease in the speed to 1 - 5 rps and a corresponding increase in pumping [4] should ensure a generation duration on the order of 500  $\mu\text{sec}$  at a scanning range  $500 \text{ cm}^{-1}$ .

#### NATURE OF CONCENTRATION QUENCHING OF THE LUMINESCENCE OF $\text{Nd}^{+++}$ IONS IN CRYSTALS

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ZhETF Pis'ma 5, No. 10, 357-360 (15 May 1967)

We present in this communication preliminary results of an investigation of the nature of concentration quenching in crystals of fluorite ( $\text{CaF}_2$ ) with  $\text{Nd}^{+++}$  added. It was shown earlier [1-3] that, besides isolated centers containing a single  $\text{Nd}^{+++}$  ion (L centers), there are produced in these crystals paired centers of two types (M and N), the concentration of which increases rapidly with increasing total  $\text{Nd}^{+++}$  concentration in the fluorite. It was shown that quenching in these crystals begins already at  $\text{Nd}^{+++}$  concentrations  $\sim 0.5 \text{ wt. } \%$ .

In the present investigation we measured the relative quantum yield  $\eta$  and the radiative lifetimes  $\tau_{\text{rad}}$  separately for single (L) and paired (M and N)  $\text{Nd}^{+++}$  centers as functions of

speed is made possible by operating in different interference orders of the Fabry-Perot cavity.

The experiments were performed on a neodymium-glass laser (Fig. 1) using samples of

KGSS-7 glass of 8 mm diameter and 80 and 120 mm long. The resonator was made up of a flat dielectric mirror 2 ( $R \sim 99\%$ ) and a total-internal reflection prism 3 rotating at speeds from 20 to 300 rps. The direction of rotation corresponded to the increase of the frequency during the process of generation. The resonator dispersion was produced by two glass prisms 4 and amounted to 4 seconds of angle per  $\text{cm}^{-1}$ . The kinetics and generation spectra were investigated in the  $9434 \text{ cm}^{-1}$  band at different pump energies and scanning rates. Depending on the level of the initial overpopulation, which was determined by the delay of the start of generation relative to the ignition of the flash lamp, we observed either continuous generation (Fig. 2a) or regular oscillations (Fig. 2b). In the latter case, generation on the order of 40  $\mu\text{sec}$  was obtained at a speed of 20 rps with the pump energy exceeding threshold by a factor of two and in a scanning range up to  $300 \text{ cm}^{-1}$ .

Figure 2c shows an oscillogram of the radiation from a laser operating in the mixed mode, when the intensity pulsations are observed against the background of a dc component. The corresponding spectrum is shown in Fig. 2d. The time scan shows that each band in the spectrum corresponds as a rule to a definite intensity burst. The interval between bands amounts on the average to  $15 - 30 \text{ cm}^{-1}$ , apparently determined by the homogeneous broadening of the  $\text{Nd}^{+++}$  ion in the glass. The band structure of the spectrum is due to the dips in the inverse population, within the limits of this value, caused by the stimulated

Fig.1. Laser with frequency scanning during generation: 1 - active medium, 2,3 - resonator reflectors, 4 - dispersive prism.

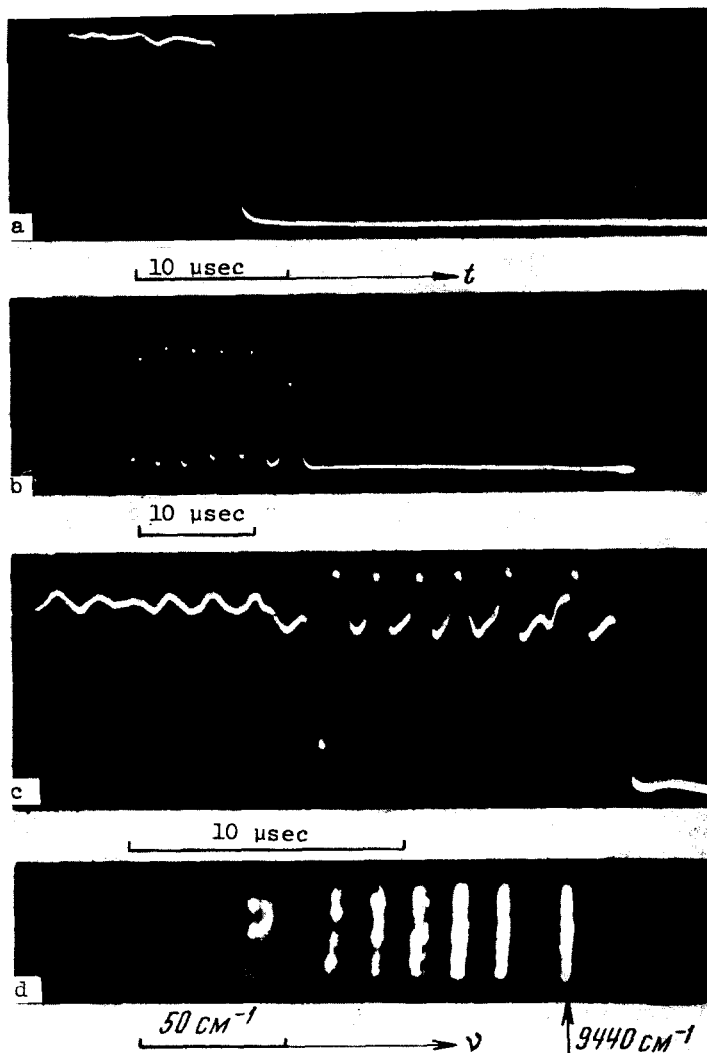
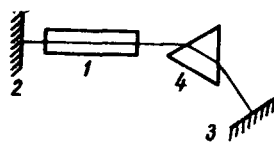


Fig.2. Oscillograms and spectrum of frequency-scanning Nd laser: a - quasi-cw mode; b - regular pulsation mode (added attenuator in front of receiver); c,d - emission oscillogram and corresponding spectrum. Prism 3 rotated at 20 rps.