

2"). The resonator consisted of external confocal dielectric mirrors ($R = 500$ mm) with 0.6% transmission. The generation was excited by a xenon flash lamp (IFP-800) in an elliptical cylindrical illuminator (semiaxes $L_1 = 16$ mm and $L_2 = 14$ mm). The stimulated emission was recorded with an FEU-28 photomultiplier with a silicon input filter. The signal from the photomultiplier was fed to an S-1-17 oscilloscope. The generation threshold was 5.6 J without allowance for the inequality of the lengths of the luminous column of the lamp and of the sample. With allowance for this inequality, the threshold was approximately 1 J. Under similar conditions, the threshold of a calcium tungstate crystal of comparable dimensions and optical quality was 2.4 J.

Figure 2 shows a photograph of the generation spectrum, obtained with a PGS-2 diffraction spectrograph. The reference was the set of an iron-arc lines in third order. The generation wavelength was 10631 \AA and the generation line width 0.2 \AA . Thus, FA:Nd^{3+} crystals have optical properties (narrow and intense luminescence line and exceptionally large level splitting) and generation parameters (low threshold and high gain) making it a promising material for continuous and pulsed lasers, and also for the investigation of the magnetic properties of the Nd^{3+} ion.

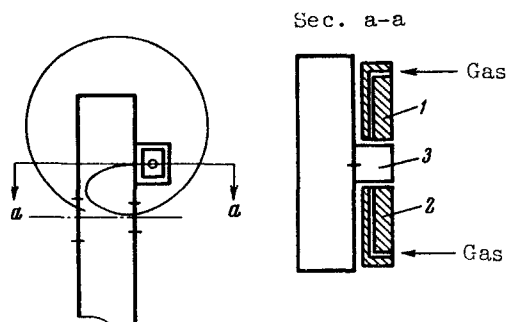
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MICROTRON WITH PLASMA INJECTOR

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In modern strong-current microtrons, the injection is by means of a scheme proposed by S. P. Kapitza, wherein the electron emitter is a thermionic cathode installed near the injection aperture of the resonator [1]. A number of limitations imposed by the thermionic cathode on the capabilities of the microtron can be eliminated in principle by using a plasma injector, which ensures large density of the currents extracted from the plasma surface, is free of incandescent parts, and has little sensitivity to the gas atmosphere in the vacuum chamber.

We started up a microtron in which the electrons were injected into the accelerating resonator from the plasma of a low-pressure gas discharge localized near the injection aperture. The injection scheme, shown in the figure, is similar to that developed earlier [1], but the thermionic cathode was replaced by a Penning discharge produced in the microtron magnetic field. The electrody system of the discharge chamber consists of two cold aluminum cathodes 1 and 2, and a copper anode made up of part 3, mounted on the cover of the resonator, and the part of the cover having the injection aperture (2.8 mm diameter). To improve the discharge-chamber characteristics, the gas (air) was admitted through



cavities in the cathodes, making it possible to realize the hollow-cathode effect in the Penning electrode system. The discharge chamber was sealed off from the high-vacuum volume by means of a polyethylene ring and communicated with this volume only through the injection aperture. The electrons were extracted from the plasma with a high-frequency field extending from the resonator cavity through the injection aperture in the discharge region.

The discharge power supply generated pulses of 30 usec duration with a repetition frequency 1 - 10 Hz. At a gas flow of 60 cm³/hr and a discharge voltage 2 kV, the discharge current was 4 A (pulsed). The pressure in the microtron chamber was 3 x 10⁻⁵ Torr. The current of the electrons accelerated to 6 MeV reached 15 mA. With the discharge sufficiently stable, instability of the accelerated current appeared with increasing pressure.

Work on the stabilization of the accelerated current is continuing.

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CONTROL OF THE FREQUENCY OF A CO₂ LASER BY A BORON TRICHLORIDE FILTER

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We present here the results of experiments aimed at controlling the frequency of a CO₂ laser by introducing in the laser resonator a cell containing BCl₃ at pressures 10⁻² - 30 Torr.

The maximum of the radiation of a CO₂ laser without a dispersive element occurs at the rotational transition P(20) of the 00⁰1 - 10⁰0 vibrational band. When the dispersive element is placed in the resonator, generation is possible at practically all the transitions of the P and R branches of the vibrational bands 00⁰1 - 10⁰0 and 00⁰1 - 02⁰0 (see, e.g., the review [1]).

The form and mutual arrangement of the absorption lines ₃ of the molecules B¹⁰Cl₃ and B¹¹Cl₃ (see Fig. 1), lying in the region of the possible CO₂-laser generation frequencies, make it possible to use BCl₃ as a selective element for the control of the frequency of this laser.

The experiments were performed on a laser with a resonator 200 cm long, made up of gold-coated spherical mirrors with curvature radius 500 cm. The diameter of the exit aperture was 0.7 cm. A cell with BCl₃, of 100 cm length and 5.5 cm diameter, was inserted in the resonator.

In the absence of BCl₃ (pressure in cell 10⁻² Torr), generation takes place at

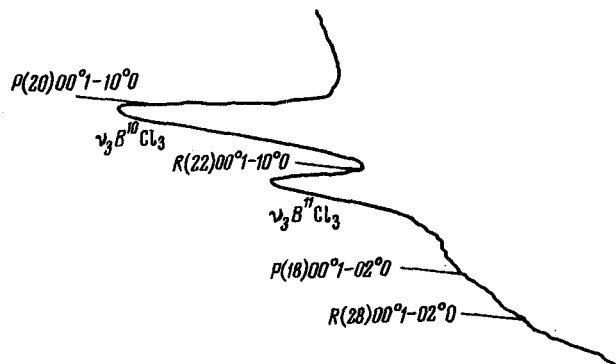


Fig. 1. Absorption spectrum of BCl₃ gas together with P and R branches of the 00⁰1 - 10⁰0 and 00⁰1 - 02⁰0 vibrational bands.