## Properties of semiconducting $Cd_xHg_{1-x}$ Te solid solutions under pressure

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We have investigated the galvanomagnetic properties of the superconducting solid solutions  $Cd_x Hg_{1-x}Te$  (x=0.2 and x=0.23) at pressures up to 14 katm. It is shown that hydrostatic compression eliminates the low-temperature anomalies, the nature of which is the same for all values of the parameter x (x < 0.23).

The change in the properties of p-type  $Cd_xHg_{1-x}Te$  solid solutions under pressure, at values of x corresponding to the semimetallic composition, was investigated in detail in<sup>[1,2]</sup> in a study of the transition to the gapless state. It was shown<sup>[2]</sup> that the galvanomagnetic phenomena and the transport phenomena are deter-

mined by the presence of acceptor levels that are superimposed on the continuous-spectrum states and split into an impurity acceptor band at impurity-center concentrations on the order of  $10^{17}$  cm<sup>-3</sup>. <sup>[3]</sup> The impurity-band conductivity constitutes in this case a noticeable (or even predominant) fraction of the total conductivity.

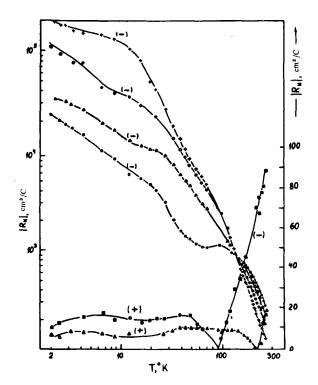


FIG. 1. Dependence of the Hall constant  $R_H$  on the temperature at different pressures for sample 5-2; +) P=0.15 katm,  $\bigcirc$ ) P=1.3 katm,  $\triangle$ ) P=1.8 katm,  $\bigcirc$ ) P=2.65 katm, left-hand scale;  $\bigcirc$ ) P=4.4 katm,  $\triangle$ ) P=8.95 katm, right-hand scale.

We report here the results of experiments on the action of hydrostatic compression on the properties of semiconducting p-type solid solutions. The samples have the following parameters at P=1 atm and  $T=4.2\,^{\circ}\mathrm{K}$ : sample 5–1)  $n=1.6\times10^{12}~\mathrm{cm^{-3}}~\mu_n=1.1\times10^6~\mathrm{cm^2/V}~\mathrm{sec},~p=1\times10^{17}~\mathrm{cm^{-3}},~\mu_p=160~\mathrm{cm^2/V}~\mathrm{sec},~E_{\mathrm{g}}=60~\mathrm{meV}$ ; sample 5–2)  $n=2.5\times10^{13}~\mathrm{cm^{-3}},~\mu_n=2\times10^5~\mathrm{cm^2/V}~\mathrm{sec},~p=3.3\times10^{16}~\mathrm{cm^{-3}},~\mu_p=270~\mathrm{cm^2/V}~\mathrm{sec},~\mathrm{and}~E_{\mathrm{g}}$ 

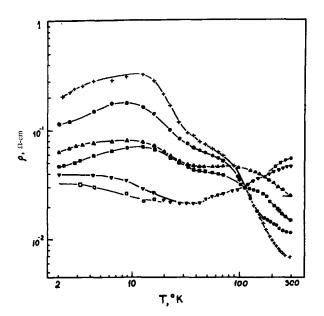


FIG. 2. Temperature dependence of the resistivity  $\rho$  of sample 5—2 at different pressures: +) P=0.15 katm, •) P=1.3 katm, •) P=1.8 katm, •) P=2.65 katm, •) P=4.4 katm, □) P=8.95 katm.

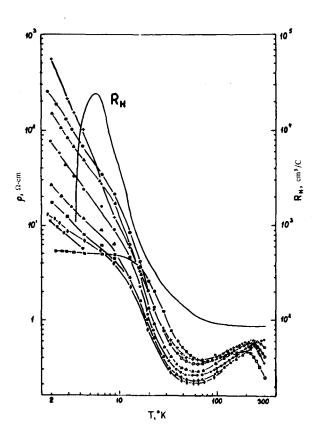


FIG. 3. Temperature dependence of the resistivity at various pressures and of the Hall constant at P=12.5 katm for sample 5-1:  $\Box$ ) P=0.15 katm, +) P=1.3 katm,  $\bigcirc$ ) P=1.8 katm,  $\triangle$ ) P=2.65 katm, •) P=4.4 katm,  $\triangle$ ) P=8.95 katm, •) P=9.7 katm, •) P=12.7 katm,  $\nabla$ ) P=14.1 katm.

= 120 meV. The parameters of the samples were calculated from the results of galvanomagnetic measurements in a weak field, and the forbidden-band width  $E_s$  was determined by optical methods, from the slopes of the  $\rho(T)$  in the high-temperature region and from the values of  $R_H$  at  $T=300\,^{\circ}{\rm K}$ ,  $^{(4)}$  and was calculated from the values of x measured with the "CAMECA" x-ray microanalyzer.  $^{[5]}$  The data obtain by different methods agree well with one another.

Figure 1 shows the temperature dependence of the Hall constant  $R_H$  at pressures  $0.15 \le p \le 8.95$  katm for sample 5-2. Just as in the case of the semimetallic samples,  $R_H$  reverses sign at a pressure  $P_{cr}$  ranging from 2.65 to 4.4 katm. The measurements were performed in a weak magnetic field,  $\mu_n H/c \ll 1$ . The transverse magnetoresistance of sample 5-2 at pressures  $P \le 2.65$  katm decreases monotonically with temperature, just as when the pressure is increased from 0.15 to 2.65 katm at a fixed temperature. The reversal of the sign of the Hall constant at  $P = P_{cr}$  is accompanied by an abrupt decrease of the transverse magnetoresistance constant  $\beta = \Delta \rho_1/\rho_0 H^2$ , from a value  $\beta = 2 \times 10^{-4}$  Oe<sup>-2</sup> at  $P = 2.65 \text{ katm to } \beta < 10^{-8} \text{ Oe}^{-2} \text{ at } P = 4.4 \text{ katm. It should}$ be noted that in a strong magnetic field (H=4 kOe) the ratio  $\Delta \rho_{1}/\rho_{0}$  reaches the anomalously large value  $\Delta \rho_{1}/$  $\rho_0$  ~60 at atmospheric pressure.

Plots of the resistivity  $\rho$  against the temperature at different pressures are shown in Fig. 2 for sample

5-2. Similar plots for sample 5-1 are shown in Fig. 3. sample 5-1 when the pressure is increased from 0.15 to 1.3 katm are still unclear. The Hall constant of sample 5-1 in a weak field has

a double inversion of the sign when the temperature is varied from 2 to 300 °K at  $P \le 2.65$  katm. When the pressure is increased to 4.4 katm and more, the Hall constant of this sample is positive at all temperatures  $T \le 300$  °K. A plot of  $R_H(T)$  for sample 5-1 at P = 12.5katm is also shown in Fig. 3.

A characteristic feature of the investigated samples is the decrease of the resistivity with pressure (at  $P \ge 1.3$  katm for 5-1) and the presence of "activation" sections on the  $\rho(T)$  and  $R_{\mu}(T)$  curves (more pronounced for sample 5-1) at low temperatures. It appears that hopping conductivity over the acceptor centers is observed in the temperature region  $2 \le T \le 10^{\circ}$ K (Fig. 3). The slopes of the  $\rho(T)$  curves at 10 < T < 25 °K are practically independent of the pressure and the activation energy of the hole conductivity corresponding to this slope is approximately 10 meV. An "activation" section is observed also on the  $R_H(T)$  curves (Fig. 3).

The causes of the abrupt jump of the resistivity of

The equality of the pressures at which the sign of the Hall constant reverse in samples 5-1 and 5-2 ( $P_{cr}$ =3-3.5 katm for both samples), which have essentially different forbidden band widths, and also the correlation with the data of [2] for semimetallic compositions. allow us to assume that the low-temperature anomalies of the galvanomagnetic properties in Cd, Hg1, Te single crystals are similar in nature for all values of the parameter x ( $x \le 0.23$ ).

<sup>&</sup>lt;sup>1</sup>T. C. Elliott, J. Melngailis, T. C. Harman, J. A. Kafalas, and W. C. Kernan, Phys. Rev. B5, 2985 (1972). <sup>2</sup>N. B. Brandt, O. N. Belousova, L. A. Bovina, V. I. Stafeev,

and Ya.G. Ponomarev, Zh. Eksp. Teor. Fiz. 66, 330 (1974) [Sov. Phys.-JETP 39, No. 1 (1974)]. <sup>3</sup>B. L. Gel'mont and M. I. D'yakonov, Zh. Eksp. Teor. Fiz.

<sup>62, 713 (1972) [</sup>Sov. Phys.-JETP 35, 377 (1972)].

<sup>&</sup>lt;sup>4</sup>C. Vérié, Phys. Stat. Sol. 17, 889 (1966).

<sup>&</sup>lt;sup>5</sup>J. L. Schmit and E. L. Stelzer, J. Appl. Phys. 40, 4865

<sup>(1969).</sup>