

## Observation of narrow $N^+(1685)$ and $N^0(1685)$ resonances in $\gamma N \rightarrow \eta\pi N$ reactions

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The observation of a narrow structure at  $W \sim 1.68$  GeV in the  $\gamma n \rightarrow \eta n$  excitation function at GRAAL, CBELSA/TAPS, LNS and A2@MAMI C [1–9] and in Compton scattering on the neutron  $\gamma n \rightarrow \gamma n$  [10] (the so-called “neutron anomaly”), and two narrow structures at  $W \sim 1.68$  GeV and  $W \sim 1.72$  GeV in Compton scattering on the proton [11], in the precise data for the  $\gamma n \rightarrow \eta n$  [12] and  $\pi^- p \rightarrow \pi^- p$  [13] reactions may signal the existence of one ( $N(1685)$ ) or two ( $N(1685)$  and  $N(1726)$ ) narrow nucleon resonances [14–19].

On the other hand, there are alternative interpretations of the “neutron anomaly” in terms of the specific interference of known wide resonances [20–24] or as the sub-threshold meson-nucleon production (cusp) [25, 26]. Although being questionable [27], the first assumption is widely discussed in literature.

The decisive identification of these experimental findings is a challenge for both theory and experiment. In the previous experiments the possible signal of  $N(1685)$  was observed in so-called “formation” reactions in which the incoming particle interacts with the target nucleon and excites resonances. If  $N(1685)$  does really exist, its signal should also be seen in multi-particle “production” reactions in which it would manifest itself as a peak in the invariant mass spectra of the final-state products. Possible reactions could be  $\gamma N \rightarrow \pi\eta N$ .

In this Letter, we report on the study of the  $\gamma p \rightarrow \pi^0\eta p$ ,  $\gamma p \rightarrow \pi^+\eta n$ ,  $\gamma n \rightarrow \pi^0\eta n$ , and  $\gamma n \rightarrow \pi^-\eta p$  reactions using the data collected at the GRAAL facility [28].

Photons from  $\eta \rightarrow 2\gamma$  and  $\pi^0 \rightarrow 2\gamma$  decays and charged pions were detected in the BGO Ball. The recoil protons and neutrons emitted at forward angles  $\theta_{\text{lab}} \leq 25^\circ$  were detected in the assembly of forward

detectors. It consisted of two planar wire chambers, a thin scintillator hodoscope and a lead-scintillator wall.

At the first step of the data analysis  $\eta$  and  $\pi^0$  mesons were identified by means of the invariant masses of two properly chosen photons. Then the cuts on the proton and  $\eta$  missing masses were applied.

At the second stage of the data analysis the cuts on the coplanarity and on the differences between the missing and invariant masses assuming two-body reactions with one real particle and one “effective” two-particle in the final state were imposed.

Given the goal of this work, only the events in the range of the energy of the incoming photon  $E_\gamma = 1.4 - 1.5$  GeV were selected for the further analysis. The lower limit of 1.4 GeV is close to the  $\gamma N \rightarrow \pi N(1685)$  threshold. The upper value 1.5 GeV is the limit of the GRAAL beam and it also allows to avoid the contribution from higher-lying resonances. To eliminate the contamination of  $\gamma N \rightarrow \eta\Delta$  events, the cuts on the invariant mass  $1.12 \leq IM(\pi N) \leq 1.22$  GeV and the missing mass  $MM(\gamma, \eta) \leq 1.22$  GeV were applied.

Fig. 1 shows the obtained  $M(\eta N)$  spectrum (the sum of all reactions under study). There is a well pronounced peak at  $\sim 1.68$  GeV. The Gaussian+3-order polynomial (signal-plus-background) fit results in the  $\chi$ -square of 23.9/23. The fit by 3-order polynomial (background) gives the  $\chi$ -square of 42.6/26. The  $\log$  likelihood ratio of these two hypotheses ( $\sqrt{2 \ln(L_{B+S}/L_B)}$ ) corresponds to the confidence level of  $4.6\sigma$ .

The extracted peak position is  $M = 1678 \pm 0.8_{\text{stat}} \pm 10_{\text{sys}}$  MeV. The systematic uncertainty in the mass position originates from the uncertainties in the calibration of the GRAAL detector and tagger.

Our results support the existence of two narrow resonances,  $N^+(1685)$  decaying, in particular, into  $\eta p$  final state, and  $N^0(1685)$  with one possible decay into  $\eta n$  (i.e. the isospin-1/2  $N(1685)$  resonance). It is unclear

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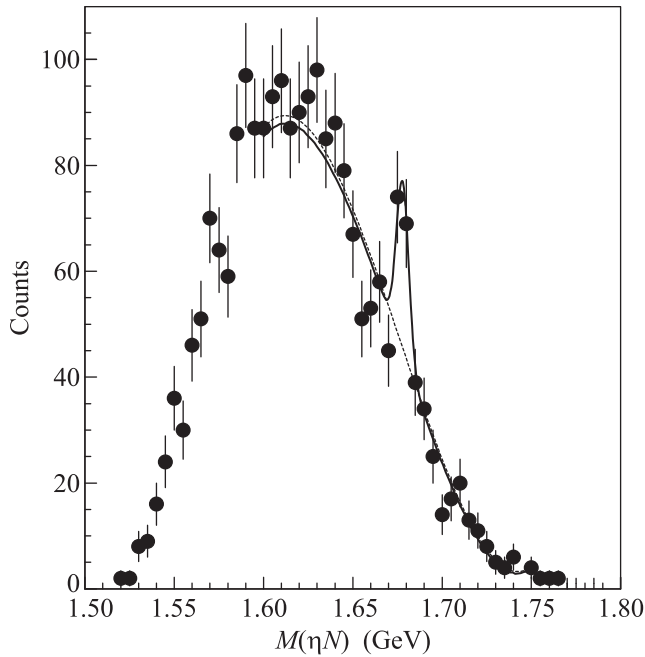


Fig. 1. Spectrum of extracted  $M(\eta N)$  mass (sum of all channels) with corrections

if the interference of known wide resonances [20–24] or the cusp effect [25, 26] – two other hypotheses under discussion – could explain these results.

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